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Invariance Principle for the Random Lorentz Gas—Beyond the Boltzmann-Grad Limit

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Dedicated to Oliver Penrose on his 91st birthday.

Abstract: We prove the invariance principle for a *random Lorentz-gas* particle in 3 dimensions under the Boltzmann-Grad limit and simultaneous diffusive scaling. That is, for the trajectory of a point-like particle moving among infinite-mass, hard-core, spherical scatterers of radius r , placed according to a Poisson point process of density ρ , in the limit $\rho \rightarrow \infty$, $r \rightarrow 0$, $\rho r^2 \rightarrow 1$ up to time scales of order $T = o(r^{-2} |\log r|^{-2})$. To our knowledge this represents the first significant progress towards solving rigorously this problem in classical nonequilibrium statistical physics, since the groundbreaking work of Gallavotti (Phys Rev 185:308–322, 1969, Nota Interna Univ di Roma 358, 1970, Statistical mechanics. A short treatise. Theoretical and mathematical physics series, Springer, Berlin, 1999), Spohn (Commun Math Phys 60:277–290, 1978, Rev Mod Phys 52:569–611, 1980) and Boldrighini–Bunimovich–Sinai (J Stat Phys 32:477–501, 1983). The novelty is that the diffusive scaling of particle trajectory and the kinetic (Boltzmann-Grad) limit are taken simultaneously. The main ingredients are a coupling of the mechanical trajectory with the Markovian random flight process, and probabilistic and geometric controls on the efficiency of this coupling. Similar results have been earlier obtained for the weak coupling limit of classical and quantum random Lorentz gas, by Komorowski–Ryzhik (Commun Math Phys 263:277–323, 2006), respectively, Erdős–Salmhofer–Yau (Acta Math 200:211–277, 2008, Commun Math Phys 271:1–53, 2007). However, the following are substantial differences between our work and these ones: (1) The physical setting is different: low density rather than weak coupling. (2) The method of approach is different: probabilistic coupling rather than analytic/perturbative. (3) Due to (2), the time scale of validity of our diffusive approximation—expressed in terms of the kinetic time scale—is much longer and fully explicit.

1. Introduction

We consider the Lorentz gas with randomly placed spherical hard core scatterers in \mathbb{R}^d . That is, place spherical balls of radius r and infinite mass centred on the points of a

Poisson point process of intensity ϱ in \mathbb{R}^d , where $r^d \varrho$ is sufficiently small so that with positive probability there is free passage out to infinity, and define $t \mapsto X^{r,\varrho}(t) \in \mathbb{R}^d$ to be the trajectory of a point particle starting with randomly oriented unit velocity, performing free flight in the complement of the scatterers and scattering elastically on them.

A major problem in mathematical statistical physics is to understand the diffusive scaling limit of the particle trajectory

$$t \mapsto \frac{X^{r,\varrho}(Tt)}{\sqrt{T}}, \quad \text{as } T \rightarrow \infty. \quad (1)$$

Indeed, the *Holy Grail* of this field of research would be to prove the invariance principle (i.e. weak convergence to a Wiener process with nondegenerate variance) for the sequence of processes in (1) in either the quenched or annealed setting (discussed in Sect. 1.1). For extensive discussion and historical background see the surveys [8, 24, 28] and the monograph [29].

The same problem in the periodic setting, when the scatterers are placed in a periodic array and randomness comes only with the initial conditions of the moving particle, is much better understood, due to the fact that in the periodic case the problem is reformulated as diffusive limit of particular additive functionals of billiards in compact domains and thus heavy artillery of hyperbolic dynamical systems theory is efficiently applicable. In order to put our results in context, we will summarise very succinctly the existing results, in Sect. 1.4.

There has been, however, no progress in the study of the *random* Lorentz gas informally described above, since the ground-breaking work of Gallavotti [15–17], Spohn [27, 28] and Boldrighini–Bunimovich–Sinai [4] where weak convergence of the process $t \mapsto X^{r,\varrho}(t)$ to a continuous time random walk $t \mapsto Y(t)$ (called Markovian flight process) was established in the Boltzmann-Grad (a.k.a. low density) limit $r \rightarrow 0$, $\varrho \rightarrow \infty$, $r^{d-1}\varrho \rightarrow 1$, in compact time intervals $t \in [0, T]$, with $T < \infty$, in the annealed [15–17, 27, 28], respectively, quenched [4] setting.

Our main result (see Theorem 2 in Sect. 1.3) proves the invariance principle in the annealed setting if we take the *Boltzmann-Grad and diffusive limits simultaneously*: $r \rightarrow 0$, $\varrho \rightarrow \infty$, $r^{d-1}\varrho \rightarrow 1$ and $T = T(r) \rightarrow \infty$. Thus while the diffusive limit (1) with fixed r and ϱ remains open, this is the first result proving convergence for times growing to infinity as $r \rightarrow 0$ in the setting of randomly placed scatterers, and hence it is a significant step towards the full resolution of the problem in the annealed setting.

1.1. The random Lorentz gas. We define now more formally the random Lorentz process. Place spherical balls of radius r and infinite mass centred on the points of a Poisson point process of intensity ϱ in \mathbb{R}^d , and define the trajectory $t \mapsto X^{r,\varrho}(t) \in \mathbb{R}^d$ of a particle moving among these scatterers as follows:

- If the origin is covered by a scatterer then $X^{r,\varrho}(t) \equiv 0$.
- If the origin is not covered by a scatterer then $t \mapsto X^{r,\varrho}(t)$ is the trajectory of a point-like particle starting from the origin with random velocity sampled uniformly from the unit sphere S^{d-1} and flying with constant speed between successive elastic collisions on any one of the fixed, infinite mass scatterers.

The randomness of the trajectory $t \mapsto X^{r,\varrho}(t)$ (when not identically 0) is due to two sources: the random placement of the scatterers and the random choice of initial velocity

of the moving particle. Otherwise, the dynamics of the moving particle is fully deterministic, governed by classical Newtonian laws. With probability 1 (with respect to both sources of randomness) the trajectory $t \mapsto X^{r,\varrho}(t)$ is well defined.

Due to elementary scaling and percolation arguments

$$\mathbf{P}(\text{the moving particle is not trapped in a compact domain}) = \vartheta_d(\varrho r^d), \quad (2)$$

where $\vartheta_d : \mathbb{R}_+ \rightarrow [0, 1]$ is a percolation probability which is (i) monotone non-increasing; (ii) continuous except for one possible jump at a positive and finite critical value $u_c = u_c(d) \in (0, \infty)$; (iii) vanishing for $u \in (u_c, \infty)$ and positive for $u \in (0, u_c)$; (iv) $\lim_{u \rightarrow 0} \vartheta_d(u) = 1$. We assume that $\varrho r^d < u_c$. In fact, in the Boltzmann-Grad limit considered in this paper (see (3) below) we will have $\varrho r^d \rightarrow 0$.

As discussed above, the Holy Grail of this field is a mathematically rigorous proof of the invariance principle of the processes (1) in either one of the following two settings.

- (Q) *Quenched limit*: For almost all (i.e. typical) realisations of the underlying Poisson point process, with averaging over the random initial velocity of the particle. In this case, it is expected that the variance of the limiting Wiener process is deterministic, not depending on the realisation of the underlying Poisson point process.
- (AQ) *Averaged-quenched* (a.k.a. *annealed*) *limit*: Averaging over the random initial velocity of the particle *and* the random placements of the scatterers.

Remarks on the Hamiltonian character of the problem: We use a probabilistic language and setting in this paper, previously much of the literature has chosen to work in the Hamiltonian setting [15–17, 27, 28]. However, we should emphasise that this probabilistic description is equivalent to the annealed setting of a Hamiltonian system: The Lorentz particle moves according to Newton’s Second Law in the potential field of spherical hard core scatterers centred in the points of a Poisson Point Process. The potential field is

$$\Phi^{r,\varrho}(x) := \sum_{q \in \omega_\varrho} \varphi((x - q)/r),$$

where ω_ϱ is the realisation of a Poisson Point Process of intensity ϱ in $\mathbb{R}^d \setminus B_{0,r}$ (that is, no scatterer within distance r from the origin) and $\varphi(x) := \infty \mathbb{1}_{\{|x| < 1\}}$ is a spherical hard-core potential. The Hamiltonian equations of motion of the Lorentz particle are *formally* written as follows

$$\dot{X}^{r,\varrho}(t) = V^{r,\varrho}(t), \quad \dot{V}^{r,\varrho}(t) = -\text{grad } \Phi^{r,\varrho}(X^{r,\varrho}(t)),$$

with initial conditions

$$X^{r,\varrho}(0) = 0, \quad V^{r,\varrho}(0) \in S^{d-1}.$$

However, since the interaction potential is *hard core*, the equations of motion are singular and should be taken with a grain of salt.

1.2. The Boltzmann-Grad limit. The Boltzmann-Grad limit is the following low (relative) density limit of the scatterer configuration:

$$r \rightarrow 0, \quad \varrho \rightarrow \infty, \quad \varrho r^{d-1} \rightarrow v_{d-1}, \quad (3)$$

where v_{d-1} is the area of the $(d-1)$ -dimensional unit disc. In this limit the expected free path length between two successive collisions will be 1. Other choices of $\lim \varrho r^{d-1} \in (0, \infty)$ are equally legitimate and would change the limit only by a time (or space) scaling factor.

It is not difficult to see that in the averaged-quenched setting and under the Boltzmann-Grad limit (3) the distribution of the first free flight length starting at any deterministic time, converges to an $EXP(1)$ and the jump in velocity after the free flight happens in a Markovian way with transition kernel

$$\mathbf{P}(v_{\text{out}} \in dv' \mid v_{\text{in}} = v) = \sigma(v, v') dv', \quad (4)$$

where dv' is the surface element on S^{d-1} and $\sigma : S^{d-1} \times S^{d-1} \rightarrow \mathbb{R}_+$ is the normalised differential cross section of a spherical hard core scatterer, computable as

$$\sigma(v, v') = \frac{1}{4v_{d-1}} |v - v'|^{3-d}. \quad (5)$$

Note that in 3-dimensions the transition probability (4) of velocity jumps is uniform. That is, the outgoing velocity v_{out} is uniformly distributed on S^2 , independently of the incoming velocity v_{in} .

It is intuitively compelling but far from easy to prove that under the Boltzmann-Grad limit (3)

$$\left\{ t \mapsto X^{r, \varrho}(t) \right\} \Rightarrow \left\{ t \mapsto Y(t) \right\}, \quad (6)$$

where the symbol \Rightarrow stands for weak convergence (of probability measures) on the space of continuous trajectories in \mathbb{R}^d , see [2]. The process $t \mapsto Y(t)$ on the right hand side is the Markovian random flight process consisting of independent free flights of $EXP(1)$ -distributed length, with Markovian velocity changes according to the scattering transition kernel (4). A formal construction of the process $t \mapsto Y(t)$ is given in Sect. 2.1. The limit (6), valid in any compact time interval $t \in [0, T]$, $T < \infty$, is rigorously established in the averaged-quenched setting in [15–17, 27, 28], and in the quenched setting in [4]. In [27] more general point processes of the scatterer positions, with sufficiently strong mixing properties are considered.

The limiting Markovian flight process $t \mapsto Y(t)$ is a continuous time random walk. Therefore, by taking a second, diffusive limit *after* the Boltzmann-Grad limit (6), Donsker's theorem (see [2]) yields indeed the invariance principle,

$$\left\{ t \mapsto T^{-1/2} Y(Tt) \right\} \Rightarrow \left\{ t \mapsto W(t) \right\}, \quad (7)$$

as $T \rightarrow \infty$, where $t \mapsto W(t)$ is the isotropic Wiener process in \mathbb{R}^d of non-degenerate variance. The variance of the limiting Wiener process W can be explicitly computed but its concrete value has no importance.

The natural question arises whether one could somehow interpolate between the double limit of taking first the Boltzmann-Grad limit (6) and then the diffusive limit (7) and the plain diffusive limit for the Lorentz process, (1). Our main result, Theorem 2

formulated in Sect. 1.3 gives a positive partial answer in dimension 3. Since our results are proved in three-dimensions from now on we formulate all statements in $d = 3$ rather than general dimension. However, in some comments we will refer to general dimension d , when appropriate.

1.3. Results. In the rest of the paper we assume $\varrho = \varrho(r) = \pi r^{-2}$ and drop the superscript ϱ from the notation of the Lorentz process.

Our results (Theorems 1 and 2 formulated below) refer to a coupling—joint realisation on the same probability space—of the Markovian random flight process $t \mapsto Y(t)$, and the quenched-averaged (annealed) Lorentz process $t \mapsto X^r(t)$. The coupling is informally described later in this section and constructed with full formal rigour in Sect. 2.2.

The first theorem states that in our coupling, up to time $T \ll r^{-1}$, the Markovian flight and Lorentz exploration processes stay together.

Theorem 1. *Let $T = T(r)$ be such that $\lim_{r \rightarrow 0} T(r) = \infty$ and $\lim_{r \rightarrow 0} rT(r) = 0$. Then*

$$\lim_{r \rightarrow 0} \mathbf{P}(\inf\{t : X^r(t) \neq Y(t)\} \leq T) = 0. \quad (8)$$

Remarks on Theorem 1: This result flashes some light on the strength of the probabilistic coupling method employed in this paper. In particular, with some elementary, purely probabilistic arguments it provides a formally stronger result than [15–17, 27] which state the weak limit (6) (which follows from (8)) for any fixed $T < \infty$. Note, however, that complementing the cited papers with *explicit error bounds* (which seems feasible) would give Theorem 1. So, Theorem 1 on its own is a complement to these fundamental results. The full strength of our method is truly exhibit in Theorem 2, our main result, which extends this result to time scales where nontrivial correlations already appear. However the proof of Theorem 1 is included as it sheds light on the structure of the proof of Theorem 2.

Theorem 2. *Let $T = T(r)$ be such that $\lim_{r \rightarrow 0} T(r) = \infty$ and $\lim_{r \rightarrow 0} r^2 |\log r|^2 T(r) = 0$. Then, for any $\delta > 0$,*

$$\lim_{r \rightarrow 0} \mathbf{P}\left(\sup_{0 \leq t \leq T} |X^r(t) - Y(t)| > \delta \sqrt{T}\right) = 0, \quad (9)$$

and hence

$$\left\{t \mapsto T^{-1/2} X^r(Tt)\right\} \Rightarrow \left\{t \mapsto W(t)\right\}, \quad (10)$$

as $r \rightarrow 0$, in the averaged-quenched sense. On the right hand side of (10) W is a standard Wiener process of variance 1 in \mathbb{R}^3 .

Indeed, the invariance principle (10) readily follows from the invariance principle for the Markovian flight process, (7), and the closeness of the two processes quantified in (9). So, it remains to prove (9). This will be the content of the larger part of this paper, Sects. 4–7.

The point of Theorem 2 is that the Boltzmann-Grad limit of scatterer configuration (3) and the diffusive scaling of the trajectory are done simultaneously, and not consecutively.

The memory effects due to recollisions and shading are controlled up to the time scale $T = T(r) = o(r^{-2} |\log r|^{-2})$.

Remarks on dimension: Our proof of Theorem 2 as it stands is valid in dimension $d = 3$ only. We give here some comments on this fact and some hints on what can/could be proved by appropriate extensions of our method. However, we stress that any of these extensions would require some extra technical efforts. In order to keep the length of this paper under a reasonable limit, we do not include these arguments and extensions.

- (1) Issues in dimension $d = 2$:
 - (a) Probabilistic estimates at the core of our proofs are valid (as stated and used) only in the transient dimensions of random walk, $d \geq 3$. This difference is implicit in the Green's function estimates of Sects. 3.3 and 5.2. Nevertheless, with extra effort and the cost of an extra logarithmic factor (of order $|\log r|$) the estimates in Sect. 3.3, used in the proof of Theorem 1, can be saved in $d = 2$, as well. Using these estimates (and relying on the Doeblin argument as hinted at in comment (2) below, Theorem 1, with a logarithmic factor, i.e. with $T(r) = o(r^{-1} |\log r|^{-1})$ turns out to be valid in $d = 2$, as well.
 - (b) A subtle geometric argument which will show up in Sects. 6.4–6.6 below, is valid only in $d \geq 3$, as well. This is unrelated to the recurrence/transience dichotomy and it is crucial in controlling the short range recollision and shadowing events, in the proof of Theorem 2.
- (2) The fact that in $d = 3$ the differential cross section of hard spherical scatterers is uniform on S^{d-1} , c.f. (4), (5), facilitates our arguments, since, in this case, the successive velocities of the random flight process $Y(t)$ form an i.i.d. sequence. In dimensions $d \neq 3$ this is not the case. However, this is a technical issue only, not of crucial importance in the argument. In dimensions $d > 3$ the differential cross section (5) satisfies Doeblin's condition $\inf_{v' \in S^{d-1}} \sigma(v, v') > 0$ and, using Doeblin's subtle trick, the sequence of successive velocities $(u_n)_{n \geq 1}$ can be broken up in random i.i.d. blocks of exponentially tight lengths. This way, the main probabilistic steps of proof can be saved. In dimension $d = 2$ Doeblin's condition does not hold directly, see (5). However, it holds for the second convolution power $\sigma^2(v, v') := \int_{S^{d-1}} \sigma(v, v'') \sigma(v'', v') dv''$ (that is, for the conditional distribution of velocity after two consecutive scatterings). In this way the sequence of successive velocities $(u_n)_{n \geq 1}$ can be broken up in random 1-dependent strongly stationary (rather than i.i.d.) blocks of exponentially tight lengths. The necessary bounds can be proved with the use of Green's function estimates for random walks with 1-dependent strongly stationary (rather than i.i.d.) steps.
- (3) Possible relation with singularity of the diffusion coefficient at $r \ll 1$ and certain limitations of our method: We state without proof the following estimates: In any dimension $d \geq 3$, there exist constants $C_k, k \geq 1$, such that

$$\mathbf{P}\left(\inf_{t > \tau_k} |Y(t) - Y(0)| < r\right) \leq \begin{cases} C_k r^k & \text{if } k \in [1, d-2], \\ C_{d-1} r^{d-1} |\log r| & \text{if } k = d-1, \\ C_k r^{d-1} & \text{if } k \in [d, \infty). \end{cases}$$

In plain words, these are bounds on the probability of the continuous time random walk $t \mapsto Y(t)$ returning to the r -neighbourhood of its starting point, *after k or more scattering events*. As we are not going to use these bounds in a technical sense we don't prove them in this paper. The proof is not hard, however. We present these bounds for the following two reasons:

- (a) The logarithmic factor in the case $k = d - 1$ seems to be related to the expected singularity of the diffusion coefficient in the presumed (but not proved) diffusive limit (1), at $r \ll 1$.
- (b) The fact that for $k \geq d$ the probability of recollision after k scattering events is of order r^{d-1} , no matter how large k , is a clear warning about a limitation of our method, as is. Indeed, beyond time scales of order $T(r) \sim r^{-d+1}$ recollision patterns of all kinds of complexities occur, preventing any attempt of breaking up the time-line into quasi-independent legs, in a rigorously controlled way, as done in our proof. In conclusion, with hard work (in particular, hard geometric estimates) in dimension $d > 3$ our proof could possibly be pushed up to time scales of order $T(r) = o(r^{-d+1} |\log r|^{-\alpha})$, with some $\alpha > 0$, but certainly not further than this. Our proof of Theorem 2 reaches essentially this limit, in $d = 3$. Going to time scales longer than r^{-d+1} would require some genuinely new idea.

Remarks on robustness of the method: Our coupling method is robust, and could be applied to a variety of other interaction potentials with only technical and not conceptual extra difficulties. However, it does not seem to be easily extendable to point processes with correlations.

- (a) Extending our methods to non-spherical hard-core scatterers would change the differential cross-section (5). As such, the sequence of successive velocities of the Y -process would not be i.i.d. but a genuine Markov chain. However there are probabilistic methods to handle such difficulties (e.g. using Doeblin's decomposition to independent blocks, as described in comment (2) above). For example, we quote the invariance principle for Ehrenfest's wind-tree model (with hyper-cube scatterers), where - since the geometry is simpler - in a *subsequent work* we prove a result analogous to Theorem 2 for times of order $T(r) = o(r^{-2} |\log r|^{-1})$, c.f. [23].
- (b) Extension to smooth potentials can be done as well, though this is somewhat trickier. In this case, besides changing the differential cross section (5) one should also deal with non-instantaneous interactions. This can be handled in the case of finite range smooth potentials. The coupled Markov process will be different: not simple flights with instantaneous velocity jumps but flights with sharp but smooth scatterings. For details of the realisation of this coupling see the forthcoming work [19] where the weak coupling limit is pushed beyond the kinetic time scale with a similar, but not identical, probabilistic coupling method.
- (c) In the construction of the exploration process - as a Markov process - it is essential, however, that the point process where the scatterers are centred be Poisson. Otherwise, the exploration process could not be realised as a Markov process and probably would be of not much use. (Recall that in [27] the Boltzmann-Grad limit (6) is proved for point processes with certain correlations allowed.) This is certainly a limitation of our method. However, spatially inhomogeneous Poisson Point Processes could be handled.

Remarks on time scales: In various works the kinetic and diffusive limiting procedures are parametrised in different ways. We chose $r \rightarrow 0$, $\varrho = r^{-1/(d-1)} \rightarrow \infty$. In order to gauge how far beyond the bare kinetic limit the diffusive limit is pushed, and to compare our time scale with existing results on weak coupling diffusive limits, cf [12, 13, 21] (see Sect. 1.4 below for some details), we should introduce the *kinetic time scale* T_{kin} . This is the space-time scale on which the kinetic limits [4, 10, 14–17, 20, 27] hold, if formulated as scaling limit of the microscopic trajectory. In our notation it is

$$T_{\text{kin}} := \varrho^{1/d} = r^{-(d-1)/d}. \quad (11)$$

This time scale is the reference to which the time scale of validity of the diffusive limit should be gauged. In terms of the microscopic space-time - where typical spacing between scatterers is of order 1 and the Lorentz particle travels with velocity of order 1 - our diffusive limit holds for time scales up to

$$T_{\text{diff}} = T_{\text{kin}} T \quad (12)$$

with

$$T = o\left(r^{-2} |\log r|^{-2}\right) = o\left(T_{\text{kin}}^{2d/(d-1)} (\log T_{\text{kin}})^{-2}\right) \stackrel{d=3}{=} o\left(T_{\text{kin}}^3 (\log T_{\text{kin}})^{-2}\right). \quad (13)$$

This is to be compared with the time scales of the similar-in-spirit classical [21], respectively, quantum [12, 13], weak coupling diffusive limits, cf. (16). See Sect. 1.4 below for some details.

The proof of Theorems 1 and 2 will be based on a coupling (that is: a joint realisation on the same probability space) of the Markovian flight process $t \mapsto Y(t)$ and the averaged-quenched realisation of the Lorentz process $t \mapsto X^r(t)$, such that the maximum distance of their positions up to time T be small order of \sqrt{T} . The Lorentz process $t \mapsto X^r(t)$ is realised as an *exploration* of the environment of scatterers. That is, as time goes on, more and more information is revealed about the position of the scatterers. As long as $X^r(t)$ traverses yet unexplored territories, it behaves just like the Markovian flight process $Y(t)$, discovering new, yet-unseen scatterers with rate 1 and scattering on them. However, unlike the Markovian flight process it has long memory, the discovered scatterers are placed forever and if the process $X^r(t)$ returns to these positions, recollisions occur. Likewise, the area swept in the past by the Lorentz exploration process $X^r(t)$ —that is: a tube of radius r around its past trajectory—is recorded as a domain where new collisions can not occur. For a formal definition of the coupling see Sect. 2.2. Let the associated velocity processes be $U(t) := \dot{Y}(t)$ and $V^r(t) := \dot{X}^r(t)$. These are almost surely piecewise constant jump processes. The coupling is realised in such a way, that

- (A) At the very beginning the two velocities coincide, $V^r(0) = U(0)$.
- (B) Occasionally, with typical frequency of order r mismatches of the two velocity processes occur. These mismatches are caused by two possible effects:
 - *Recollisions* of the Lorentz exploration process with a scatterer placed in the past. This causes a collision event when $V^r(t)$ changes while $U(t)$ does not.
 - Scatterings of the Markovian flight process $Y(t)$ in a moment when the Lorentz exploration process is in the explored tube, where it can not encounter a not-yet-seen new scatterer. In these moments the process $U(t)$ has a jump discontinuity, while the process $V^r(t)$ stays unchanged. We will call these events *shadowed scatterings* of the Markovian flight process.
- (C) However, shortly after the mismatch events described in item (B) above, a new jointly realised scattering event of the two processes occurs, recoupling the two velocity processes to identical values. These recouplings occur typically at an $EXP(1)$ -distributed time after the mismatches.

Summarising: The coupled velocity processes $t \mapsto (U(t), V^r(t))$ are realised in such a way that they assume the same values except for typical time intervals of length of order 1, separated by typical intervals of lengths of order r^{-1} . Other, more complicated mismatches of the two processes occur only at time scales of order $r^{-2} |\log r|^{-2}$. If the probability of all mismatches, and the separation associated to those that do occur, can be controlled (this will be the content of the proof) then the following holds:

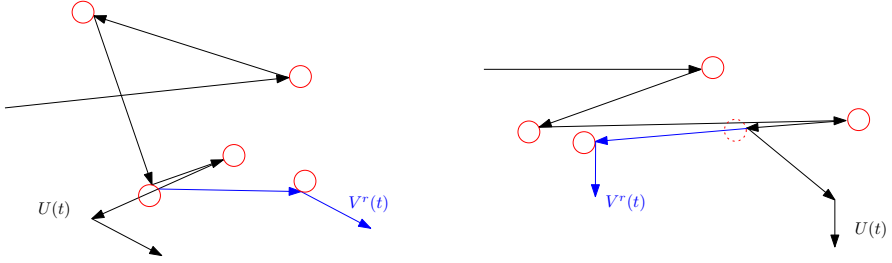


Fig. 1. The above image shows a recollision (left) and a shadowing event (right). Note that after each event U and V^r are no longer coupled. However at the next scattering, if possible, the velocities are recoupled. On the right-hand-side the virtual scatterer drawn in dotted line is shadowed. That is: it is physically not present in the mechanical trajectory

Up to $T = T(r) = o(r^{-1})$, with high probability there is no mismatch whatsoever between $U(t)$ and $V^r(t)$. That is,

$$\lim_{r \rightarrow 0} \mathbf{P}(\inf\{t : V^r(t) \neq U(t)\} < T) = \lim_{r \rightarrow 0} \mathbf{P}(\inf\{t : X^r(t) \neq Y(t)\} < T) = 0. \quad (14)$$

In particular, the invariance principle (10) also follows, with $T = T(r) = o(r^{-1})$, rather than $T = T(r) = o(r^{-2} |\log r|^{-2})$. As a by-product of this argument a new and handier proof of the theorem (6) of [15–17, 27, 28] also drops out.

Going up to $T = T(r) = o(r^{-2} |\log r|^{-2})$ needs more argument. The ideas exposed in the outline (A), (B), (C) above lead to the following chain of bounds:

$$\begin{aligned} \max_{0 \leq t \leq 1} \left| \frac{X^r(Tt)}{\sqrt{T}} - \frac{Y(Tt)}{\sqrt{T}} \right| &= \frac{1}{\sqrt{T}} \max_{0 \leq t \leq 1} \left| \int_0^{Tt} (V^r(s) - U(s)) ds \right| \\ &\leq \frac{1}{\sqrt{T}} \int_0^T |V^r(s) - U(s)| ds \asymp \frac{1}{\sqrt{T}} Tr = \sqrt{T}r. \end{aligned}$$

In the \asymp step we use the arguments (B) and (C). Finally, choosing in the end $T = T(r) = o(r^{-2})$ we obtain a tightly close coupling of the diffusively scaled processes $t \mapsto X^r(Tt)/\sqrt{T}$ and $t \mapsto Y(Tt)/\sqrt{T}$, (9), and hence the invariance principle (10), for this longer time scale. This hand-waving argument should, however, be taken with a grain of salt: it does not show the logarithmic factor, which arises in the fine-tuning.

1.4. Summary of related work. In order to put our results in context we succinctly summarise the related most important results in the mathematically rigorous treatment of diffusion in the Lorentz gas. As Hendrik Lorentz’s seminal paper [22]—where he proposes the periodic setting of what we call today the Lorentz gas for modelling diffusion and transport in solids—was published in 1905, and due to the large amount of work done in this field, we can not strive for exhaustion, and mention only a (possibly subjective) selection of the mathematically rigorous results. For more comprehensive historical overview we refer the reader to the survey papers [8, 24, 28] and the monograph [29].

Scaling limit of the periodic Lorentz gas As already mentioned, diffusion in the periodic setting is much better understood than in the random setting. This is due to the fact that diffusion in the periodic Lorentz gas can be reduced to the study of limit theorems of

some particular additive functionals of billiard flows in compact domains. Heavy tools of hyperbolic dynamics provide the technical arsenal for the study of these problems.

The first breakthrough was the fully rigorous proof, by Bunimovich and Sinai [5], of the invariance principle (diffusive scaling limit) for the Lorentz particle trajectory in a two-dimensional periodic array of spherical scatterers with finite horizon. (Finite horizon means that the length of the straight path segments not intersecting a scatterer is bounded from above.) This result was extended by Chernov [7], to higher dimensions, under a still-not-proved technical assumption on singularities of the corresponding billiard flow.

In the case of infinite horizon (e.g. the plain \mathbb{Z}^d arrangement of the spherical scatterers of diameter less than the lattice spacing) the free flight distribution of a particle flying in a uniformly sampled random direction has a heavy tail which causes a different type of long time behaviour of the particle displacement. The arguments of Bleher [3] indicated that in the two-dimensional case super-diffusive scaling of order $\sqrt{t \log t}$ is expected. For the Lorentz-particle displacement in the 2-dimensional periodic case with infinite horizon, a central limit theorem with this anomalous scaling was proved with full rigour by Varjú and Szász [30] and Dolgopyat and Chernov [9]. The periodic infinite horizon case in dimensions $d \geq 3$ remains open.

Boltzmann-Grad limit of the periodic Lorentz gas The Boltzmann-Grad limit in the periodic case means spherical scatterers of radii $r \ll 1$ placed on the points of the hypercubic lattice $r^{(d-1)/d} \mathbb{Z}^d$. The particle starts with random initial position and velocity sampled uniformly and collides elastically on the scatterers. For a full exposition of the long and complex history of this problem we quote the surveys [18,24] and recall only the final, definitive results.

In Caglioti–Golse [6] and Marklof–Strömbergsson [25] it is proved that in the Boltzmann-Grad limit the trajectory of the Lorentz particle in any compact time interval $t \in [0, T]$ with $T < \infty$ fixed, converges weakly to a non-Markovian flight process which has, however, a complete description in terms of a Markov chain of the successive collision impact parameters and, conditionally on this random sequence, independent flight lengths. (For a full description in these terms see [26].) As a second limit, the invariance principle is proved for this non-Markovian random flight process, with superdiffusive scaling $\sqrt{t \log t}$, in Marklof–Tóth [26]. Note that in this case the second limit doesn't just drop out from Donsker's theorem as it did in the random scatterer setting. The results of [6] are valid in $d = 2$ while those of [25] and [26] in arbitrary dimension.

Interpolating between the plain scaling limit in the infinite horizon case (open in $d \geq 3$) and the kinetic limit, by simultaneously taking the Boltzmann-Grad limit *and* scaling the trajectory by $\sqrt{T \log T}$, where $T = T(r) \rightarrow \infty$ with some rate, would be the problem analogous to our Theorem 1 or Theorem 2. This is widely open.

The weak coupling limit The weak coupling is physically a different limiting procedure for obtaining diffusion of moving particle among fixed scatterers. In conformity with the usual notation of the weak coupling literature we will use the scaling parameter $\varepsilon \rightarrow 0$. Infinite mass fixed scatterers are again placed on the points of a Poisson point process of density $\varrho = \varepsilon^{-d}$ in \mathbb{R}^d . However, now it is assumed that the compactly supported and spherically symmetric scattering potential \mathcal{U} of radius $r = \varepsilon$, centred at the scatterer positions, is smooth and bounded rather than hard core. Note that $\varrho = \varepsilon^{-d}$, $r = \varepsilon$ means just a linear spatial scaling by a factor ε . In this limit, rather than scaling down excessively the radius of support, the strength of the potential is scaled. Newton's equations of motion for the kinetically scaled particle are

$$\dot{X}^\varepsilon(t) = V^\varepsilon(t), \quad \dot{V}^\varepsilon(t) = -\nabla U^\varepsilon(X^\varepsilon(t))$$

in the potential field

$$U^\varepsilon(x) = \sum_{q \in \omega} \varepsilon^{1/2} \mathcal{U}(\varepsilon^{-1}(x - q)),$$

where ω is the realisation of the Poisson point process of intensity $\varrho = \varepsilon^{-d}$.

From the work of Kesten and Papanicolaou [20] it follows that

$$V^\varepsilon(t) \Rightarrow \mathcal{V}(t), \quad X^\varepsilon(t) \Rightarrow \mathcal{X}(t) := \int_0^t \mathcal{V}(s) ds, \quad (15)$$

where the limiting velocity process $\mathcal{V}(t)$ is a homogeneous diffusion (i.e. Brownian motion) on the surface of S^{d-1} and the weak convergence is meant in the space of continuous trajectories endowed with uniform topology on compact time intervals, cf [2]. See also the survey [28]. Taking a second, diffusive limit, $T^{-1/2} \mathcal{X}(Tt) \rightarrow W(t)$, the displacement process converges to Brownian motion, as $T \rightarrow \infty$.

The simultaneous kinetic *and* diffusive limit in this context is done by Komorowski and Ryzhik in [21] where it is proved that in dimension $d \geq 3$, up to time scales

$$T = T(\varepsilon) = \varepsilon^{-\kappa} = T_{\text{kin}}^\kappa, \quad \kappa \in (0, \kappa_0), \quad \kappa_0 > 0, \quad (16)$$

the diffusive limit

$$T^{-1/2} X^\varepsilon(Tt) \Rightarrow W(t) \quad (17)$$

holds. In (16) κ_0 is small (possibly, very small) and positive, its numerical value is not specified and difficult to determine from the various technical estimates.

To our knowledge this was the first case when diffusive limit was rigorously established beyond the kinetic time scale in a context which includes the random Lorentz gas. We also note that the results in [20, 21] are formulated in more general context of spatially ergodic random potential fields with regularity conditions assumed. This covers weak coupling of the random Lorentz gas as particular case. Our main Theorem 2 should be compared with this result. In particular, the time scale of validity of the diffusive limit (16) is to be compared with the time scale (13) up to which our Theorem 2 is valid.

In the forthcoming work [19] the diffusive limit under weak coupling (17) is proved with probabilistic coupling method somewhat similar but not identical to the present one, for time scales $T(\varepsilon) = o(\varepsilon^{-d+2} |\log \varepsilon|^{-\alpha})$ in any $d \geq 3$, with some $\alpha < \infty$, improving thus considerably the result of [21].

The quantum Lorentz gas The quantum versions of the weak coupling and low density limits for the random Lorentz gas were considered in Erdős-Yau [14], respectively, Eng-Erdős [10], where the long time evolution of a quantum particle interacting with a random potential is studied. It is proved that the phase space density of the quantum evolution converges weakly to the solution of the linear Boltzmann (or, Langevin) equation, with diffusive, respectively, hopping scattering kernels. These results are the quantum analogues of the classical (i.e. non-quantum) kinetic limits of [20] (for weak coupling), respectively, [15–17, 27, 28] (for low density).

In the weak coupling setup the simultaneous kinetic *and* diffusive scaling limit, formally analogous to [21] was done by Erdős-Salmhofer-Yau [12, 13] where it is proved

that under a scaling limit similar to (16), (17) the time evolution of the spatial density of the quantum particle weakly coupled with the fixed scatterers converges to the solution of the heat equation. In this case the numerical value of the upper bound on the scaling exponent κ is specified in $d = 3$ as $\kappa_0 = 1/370$ (see Theorem 2.2 in [12]).

For a comprehensive survey of the kinetic and kinetic-diffusive limits in the quantum case see also [11].

Miscellaneous Looking into the future: Liverani investigates the periodic Lorentz gas with finite horizon with local random perturbations in the cells of periodicity: a basic periodic structure with spherical scatterers centred on \mathbb{Z}^d with extra scatterers placed randomly and independently within the cells of periodicity, [1]. This is an interesting mixture of the periodic and random settings which could succumb to a mixture of dynamical and probabilistic methods, so-called deterministic walks in random environment.

1.5. Structure of the paper. The rest of the paper is devoted to the rigorous statement and proof of the arguments exposed in (A), (B), (C) above. Its overall structure is as follows:

- *Section 2:* We construct the Markovian flight and Lorentz exploration processes and thus lay out the coupling argument which is essential moving forward. Moreover, we will also introduce an auxiliary process, Z , a short-sighted or forgetful version of X which somehow interpolates between the processes Y and X .
- *Section 3:* We prove Theorem 1. We go through the proof of this statement as it is both informative for the dynamics, and the proof of Theorem 2 in its full strength will follow partially similar lines, however with substantial differences.

Sections 4–7 are fully devoted to the proof of Theorem 2, as follows:

- *Section 4:* We break up the process Z into independent legs of exponentially tight lengths. From here we state two propositions which are central to the proof. They state that
 - (i) with high probability the process X does not differ from Z in each leg;
 - (ii) with high probability, the different legs of the process Z do not interact (up to times of our time scales).
- *Section 5:* We prove the proposition concerning interactions between legs.
- *Section 6:* We prove the proposition concerning coincidence, with high probability, of the processes X and Z within a single leg. This section is longer than the others, due to the subtle geometric arguments and estimates needed in this proof.
- *Section 7:* We finish off the proof of Theorem 2.

2. Construction

2.1. Ingredients and the Markovian flight process. Let $\xi_j \in \mathbb{R}_+$ and $u_j \in \mathbb{R}^3$, $j = -2, -1, 0, 1, 2, \dots$, be completely independent random variables (defined on an unspecified probability space $(\Omega, \mathcal{F}, \mathbf{P})$) with distributions:

$$\xi_j \sim EXP(1), \quad u_j \sim UNI(S^2), \quad (18)$$

and let

$$y_j := \xi_j u_j \in \mathbb{R}^3. \quad (19)$$

For later use we also introduce the sequence of indicators

$$\epsilon_j := \mathbb{1}\{\xi_j < 1\}, \quad (20)$$

and the corresponding conditional exponential distributions $EXP(1|1) := \text{distrib}(\xi | \epsilon = 1)$, respectively, $EXP(1|0) = \text{distrib}(\xi | \epsilon = 0)$, with distribution densities

$$(e - 1)^{-1} e^{1-x} \mathbb{1}\{0 \leq x < 1\}, \quad \text{respectively,} \quad e^{1-x} \mathbb{1}\{1 \leq x < \infty\}.$$

We will also use the notation $\underline{\epsilon} := (\epsilon_j)_{j \geq 0}$ and call the sequence $\underline{\epsilon}$ the *signature* of the i.i.d. $EXP(1)$ -sequence $(\xi_j)_{j \geq 0}$.

The variables ξ_j and u_j will be, respectively, the consecutive *flight length/flight times* and *flight velocities* of the Markovian flight process $t \mapsto Y(t) \in \mathbb{R}^3$ defined below.

Denote, for $n \in \mathbb{Z}_+$, $t \in \mathbb{R}_+$,

$$\tau_n := \sum_{j=1}^n \xi_j, \quad \nu_t := \max\{n : \tau_n \leq t\}, \quad \{t\} := t - \tau_{\nu_t}. \quad (21)$$

That is: τ_n denotes the consecutive scattering times of the flight process, ν_t is the number of scattering events of the flight process Y occurring in the time interval $(0, t]$, and $\{t\}$ is the length of the last free flight before time t .

Finally let

$$Y_n := \sum_{j=1}^n \xi_j u_j = \sum_{j=1}^n y_j, \quad Y(t) := Y_{\nu_t} + \{t\} u_{\nu_t+1}.$$

We shall refer to the process $t \mapsto Y(t)$ as the Markovian flight process. This will be our fundamental probabilistic object. All variables and processes will be defined in terms of this process, and adapted to the natural continuous time filtration $(\mathcal{F}_t)_{t \geq 0}$ of the flight process:

$$\mathcal{F}_t := \sigma(u_0, (Y(s))_{0 \leq s \leq t}).$$

Note that the processes $n \mapsto Y_n$, $t \mapsto Y(t)$ and their respective natural filtrations $(\mathcal{F}_n)_{n \geq 0}$, $(\mathcal{F}_t)_{t \geq 0}$, do not depend on the parameter r .

We also define, for later use, the *virtual scatterers* of the flight process $t \mapsto Y(t)$. Let

$$Y'_k := Y_k + r \frac{u_k - u_{k+1}}{|u_k - u_{k+1}|} = Y_k + r \frac{\dot{Y}(\tau_k^-) - \dot{Y}(\tau_k^+)}{|\dot{Y}(\tau_k^-) - \dot{Y}(\tau_k^+)|}, \quad k \geq 0,$$

$$\mathcal{S}_n^Y := \{Y'_k \in \mathbb{R}^3 : 0 \leq k \leq n\}, \quad n \geq 0.$$

Here and throughout the paper we use the notation $f(t^\pm) := \lim_{\varepsilon \downarrow 0} f(t \pm \varepsilon)$.

The points $Y'_n \in \mathbb{R}^3$ are the centres of virtual spherical scatterers of radius r which *would have caused* the n th scattering event of the flight process. They do not have any influence on the further trajectory of the flight process Y , but will play role in the forthcoming couplings.

2.2. The Lorentz exploration process. Let $r > 0$, and $\varrho = \varrho(r) = \pi r^{-2}$. We define the Lorentz exploration process $t \mapsto X(t) = X^r(t) \in \mathbb{R}^3$, coupled with the flight process $t \mapsto Y(t)$, adapted to the filtration $(\mathcal{F}_t)_{t \geq 0}$. The process $t \mapsto X(t)$ and all upcoming random variables related to it *do depend* on the choice of the parameter r (and ϱ), but from now on we will suppress explicit notation of dependence upon these parameters.

The construction goes inductively, on the successive time intervals $[\tau_{n-1}, \tau_n]$, $n = 1, 2, \dots$. Start with [Step 1] and then iterate indefinitely [Step 2] and [Step 3] below.

[Step 1] Start with

$$X(0) = X_0 = 0, \quad V(0^+) = u_1, \quad X'_0 := r \frac{u_0 - u_1}{|u_0 - u_1|}, \quad \mathcal{S}_0^X = \{X'_0\}.$$

Note that the trajectory of the exploration process X begins with a collision at time $t = 0$. This is not exactly as described previously but is of no consequence and aids the later exposition. Go to [Step 2].

Step 2 This step starts with given $X(\tau_{n-1}) = X_{n-1} \in \mathbb{R}^3$, $V(\tau_{n-1}^+) \in S^2$ and $\mathcal{S}_{n-1}^X = \{X'_k : 0 \leq k \leq n-1\} \subset \mathbb{R}^3 \cup \{\star\}$, where

- \star is a fictitious point at infinity, with $\inf_{x \in \mathbb{R}^3} |x - \star| = \infty$, introduced for bookkeeping reasons;
- $|X_{n-1} - X'_k| \in (r, \infty]$ for $0 \leq k < n-1$, and $|X_{n-1} - X'_{n-1}| \in \{r, \infty\}$.
(Note that, due to absolute continuity of the flight time distribution $|X_l - X'_k| \neq r$, for $k \neq l$, with probability 1.)

The trajectory $t \mapsto X(t)$, $t \in [\tau_{n-1}, \tau_n]$, is defined as free motion with elastic collisions on fixed spherical scatterers of radius r centred at the points in \mathcal{S}_{n-1}^X . At the end of this time interval the position and velocity of the Lorentz exploration process are $X(\tau_n) =: X_n$, respectively, $V(\tau_n^-)$. Go to [Step 3].

Step 3 Let

$$X''_n := X_n + r \frac{V(\tau_n^-) - u_{n+1}}{|V(\tau_n^-) - u_{n+1}|}, \quad d_n := \min_{0 \leq s < \tau_n} |X(s) - X''_n|.$$

Note that $d_n \leq r$.

- If $d_n < r$ then let $X'_n := \star$, and $V(\tau_n^+) = V(\tau_n^-)$.
- If $d_n = r$ then let $X'_n := X''_n$, and $V(\tau_n^+) = u_{n+1}$.

Set $\mathcal{S}_n^X = \mathcal{S}_{n-1}^X \cup \{X'_n\}$. Go back to [Step 2].

The process $t \mapsto X(t)$ is indeed adapted to the filtration $(\mathcal{F}_t)_{0 \leq t < \infty}$ and indeed has the averaged-quenched distribution of the Lorentz process. This follows from the fact that the scatterers of the Lorentz process are centred on a Poisson Point Process and thus when sweeping not-yet-seen areas no information from the past interferes.

Our notation is fully consistent with the one used for the Markovian process Y : $X_n := X(\tau_n)$ and

$$X'_k := \begin{cases} X_k + r \frac{\dot{X}(\tau_k^-) - \dot{X}(\tau_k^+)}{|\dot{X}(\tau_k^-) - \dot{X}(\tau_k^+)|} & \text{if } \dot{X}(\tau_k^-) \neq \dot{X}(\tau_k^+), \\ \star & \text{if } \dot{X}(\tau_k^-) = \dot{X}(\tau_k^+), \end{cases} \quad k \geq 0,$$

$$\mathcal{S}_n^X := \{X'_k \in \mathbb{R}^3 \cup \{\star\} : 0 \leq k \leq n\}, \quad n \geq 0.$$

The second alternative above happens when the scattering event offered by the Poisson flaw is suppressed due to shadowing.

2.3. Mechanical consistency and compatibility of piece-wise linear trajectories in \mathbb{R}^3 . The key notion in the exploration construction of Sect. 2.2 was mechanical r -consistency, and r -compatibility of finite segments of piece-wise linear trajectories in \mathbb{R}^3 , which we are going to formalise now, for later reference. We will apply the notion of r -consistency/compatibility to several different process in what follows. Thus we let $\{t \mapsto \mathcal{Z}(t)\}$ denote any one of the aforementioned/forthcoming piece-wise linear processes.

Thus, let

$$n \in \mathbb{N}, \quad \tau_0 \in \mathbb{R}, \quad \mathcal{Z}_0 \in \mathbb{R}^3, \quad v_0, \dots, v_{n+1} \in S^2 \quad t_1, \dots, t_n \in \mathbb{R}_+,$$

be given and define for $j = 0, \dots, n-1$,

$$\tau_j := \tau_0 + \sum_{k=1}^j t_k, \quad \mathcal{Z}_j := \mathcal{Z}_0 + \sum_{k=1}^j t_k v_k, \quad \mathcal{Z}'_j := \begin{cases} \mathcal{Z}_j + r \frac{v_j - v_{j+1}}{|v_j - v_{j+1}|} & \text{if } v_j \neq v_{j+1}, \\ \star & \text{if } v_j = v_{j+1}, \end{cases}$$

and for $t \in [\tau_j, \tau_{j+1}]$, $j = 0, \dots, n$,

$$\mathcal{Z}(t) := \mathcal{Z}_j + (t - \tau_j)v_{j+1}.$$

We call the piece-wise linear trajectory $(\mathcal{Z}(t) : \tau_0^- < t < \tau_n^+)$ mechanically r -consistent or r -inconsistent, if

$$\min_{\tau_0 \leq t \leq \tau_n} \min_{0 \leq j < n} |\mathcal{Z}(t) - \mathcal{Z}'_j| = r, \quad \text{respectively,} \quad \min_{\tau_0 \leq t \leq \tau_n} \min_{0 \leq j < n} |\mathcal{Z}(t) - \mathcal{Z}'_j| < r. \quad (22)$$

Note, that by formal definition the minimum distance on the left hand side can not be strictly larger than r .

Given two finite pieces of mechanically r -consistent trajectories $(\mathcal{Z}_a(t) : \tau_{a,0}^- < t < \tau_{a,n_a}^+)$ and $(\mathcal{Z}_b(t) : \tau_{b,0}^- < t < \tau_{b,n_b}^+)$, defined over non-overlapping time intervals: $[\tau_{a,0}, \tau_{a,n_a}] \cap [\tau_{b,0}, \tau_{b,n_b}] = \emptyset$, with $\tau_{a,n_a} \leq \tau_{b,0}$, we will call them mechanically r -compatible or r -incompatible if

$$\begin{aligned} \min\{ \min_{\tau_{a,0} \leq t \leq \tau_{a,n_a}} \min_{0 \leq j \leq n_b} |\mathcal{Z}_a(t) - \mathcal{Z}'_{b,j}|, \min_{\tau_{b,0} \leq t \leq \tau_{b,n_b}} \min_{0 \leq j < n_a} |\mathcal{Z}_b(t) - \mathcal{Z}'_{a,j}| \} &\geq r, \\ \min\{ \min_{\tau_{a,0} \leq t \leq \tau_{a,n_a}} \min_{0 \leq j \leq n_b} |\mathcal{Z}_a(t) - \mathcal{Z}'_{b,j}|, \min_{\tau_{b,0} \leq t \leq \tau_{b,n_b}} \min_{0 \leq j < n_a} |\mathcal{Z}_b(t) - \mathcal{Z}'_{a,j}| \} &< r, \end{aligned} \quad (23)$$

respectively.

It is obvious that given a mechanically r -consistent trajectory, any non-overlapping parts of it are pairwise mechanically r -compatible, and given a finite number of non-overlapping mechanically r -consistent pieces of trajectories which are also pair-wise mechanically r -compatible their concatenation (in the most natural way) is mechanically r -consistent.

2.4. An auxiliary process. It will be convenient to introduce a third, auxiliary process $t \mapsto Z(t) \in \mathbb{R}^3$, and consider the joint realisation of all three processes $t \mapsto (Y(t), X(t), Z(t))$ on the same probability space. This construction will not be needed until Sect. 4, but this is the optimal logical point to introduce it. The reader may safely skip to Sect. 3 and come back here before turning to Sect. 4.

The process $t \mapsto Z(t)$ will be a *short-sighted* version of the true physical process $t \mapsto X(t)$ in the sense that in its construction only memory effects by the last seen scatterers are taken into account. That is: only *direct recollisions* with the last seen scatterer and *direct shadowings* by the last straight flight segment are incorporated, disregarding more complex memory effects. Later in the paper the following two basic facts will be shown:

- (a) Up to times $T = T(r) = o(r^{-2} |\log r|^{-2})$ the trajectories of the short-sighted process $Z(t)$ and the true physical process $X(t)$ coincide. This is the main part of the proof, filling Sects. 5 and 6 and concluded in Lemma 9.
- (b) The short-sighted process $Z(t)$ and the Markovian process $Y(t)$ stay sufficiently close together with probability tending to 1 (as $r \rightarrow 0$). This is the content of Lemma 10, its proof is relatively simple.

Based on these two conclusions, the invariance principle (7) can be transferred to the true physical process $X(t)$, thus yielding the invariance principle (10).

Define the following indicator variables:

$$\begin{aligned}\hat{\eta}_j &= \hat{\eta}(y_{j-2}, y_{j-1}, y_j) := \mathbb{1} \left\{ |y_{j-1}| < 1 \text{ and } \min_{0 \leq t \leq \xi_{j-2}} \left| y_{j-1} + r \frac{u_{j-1} - u_j}{|u_{j-1} - u_j|} + t u_{j-2} \right| < r \right\}, \\ \tilde{\eta}_j &= \tilde{\eta}(y_{j-2}, y_{j-1}, y_j) := \mathbb{1} \left\{ |y_{j-1}| < 1 \text{ and } \min_{0 \leq t \leq \xi_j} \left| y_{j-1} + r \frac{u_{j-1} - u_{j-2}}{|u_{j-1} - u_{j-2}|} + t u_j \right| < r \right\}, \\ \eta_j &:= \max\{\hat{\eta}_j, \tilde{\eta}_j\}.\end{aligned}\tag{24}$$

Before constructing the auxiliary process $t \mapsto Z(t)$ we prove the following

Lemma 1. *There exists a constant $C < \infty$ such that for any sequence of signatures $\epsilon = (\epsilon_j)_{j \geq 1}$ the following bounds hold*

$$\mathbf{E}(\eta_j \mid \epsilon) \leq Cr, \tag{25}$$

$$\mathbf{E}(\eta_j \eta_k \mid \epsilon) \leq \begin{cases} Cr^2 |\log r| & \text{if } |j - k| = 1, \\ Cr^2 & \text{if } |j - k| > 1. \end{cases} \tag{26}$$

Proof of Lemma 1. Define the following auxiliary, and simpler, indicators:

$$\hat{\eta}'_j := \mathbb{1} \left\{ \angle(-u_{j-1}, u_{j-2}) < \frac{2r}{\xi_{j-1}} \right\}, \quad \tilde{\eta}'_j := \mathbb{1} \left\{ \angle(-u_{j-1}, u_j) < \frac{2r}{\xi_{j-1}} \right\}.$$

Here, and in the rest of the paper we use the notation

$$\angle : S^2 \times S^2 \rightarrow [0, \pi], \quad \angle(u, v) := \arccos(u \cdot v).$$

Then, clearly,

$$\tilde{\eta}_j \leq \tilde{\eta}'_j, \quad \hat{\eta}_j \leq \hat{\eta}'_j.$$

It is straightforward that the indicators $(\tilde{\eta}_j : 1 \leq j < \infty)$, and likewise, the indicators $(\tilde{\eta}_j^* : 1 \leq j < \infty)$, are independent among themselves and one-dependent across the two sequences. This holds even if conditioned on the sequence of signatures $\underline{\epsilon}$.

Therefore, the following simple computations prove the claim of the lemma.

$$\mathbf{E}(\tilde{\eta}_j \mid \underline{\epsilon}) \leq Cr^2 \int_0^\infty e^{-y} \min\{y^{-2}, r^{-2}\} dy \leq Cr,$$

$$\mathbf{E}(\tilde{\eta}_j^* \mid \underline{\epsilon}) \leq Cr^2 \int_0^\infty e^{-y} \min\{y^{-2}, r^{-2}\} dy \leq Cr,$$

$$\mathbf{E}(\tilde{\eta}_{j+1} \tilde{\eta}_j \mid \underline{\epsilon}) \leq Cr^2 \int_0^\infty \int_0^\infty e^{-y} e^{-z} \min\{y^{-2}, z^{-2}, r^{-2}\} dy dz \leq Cr^2 |\log r|.$$

We omit the elementary computational details. \square

Lemma 1 assures that, as $r \rightarrow 0$, with probability tending to 1, up to time of order $T = T(r) = o(r^{-2} |\log r|^{-1})$ it will not occur that two neighbouring or next-neighbouring η -s happen to take the value 1 which would obscure the following construction.

The algorithmic definition of the process $t \mapsto Z(t)$, in terms of and adapted to the natural filtration of the flight process $t \mapsto Y(t)$, goes as follows. The process $t \mapsto Z(t)$ is constructed on the successive intervals $[\tau_{j-1}, \tau_j]$, $j = 1, 2, \dots$, as follows:

◦ (No interference with the past.)

If $\eta_j = 0$ then for $\tau_{j-1} \leq t \leq \tau_j$, $Z(t) = Z(\tau_{j-1}) + \{t\}u_j$.

◦ (Direct shadowing.)

If $\hat{\eta}_j = 1$, then for $\tau_{j-1} \leq t \leq \tau_j$, $Z(t) = Z(\tau_{j-1}) + \{t\}u_{j-1}$.

◦ (Direct recollision with the last seen scatterer.)

If $\hat{\eta}_j = 0$ and $\tilde{\eta}_j = 1$ then, in the time interval $\tau_{j-1} \leq t \leq \tau_j$ the trajectory $t \mapsto Z(t)$ is defined as that of a mechanical particle starting with initial position $Z(\tau_{j-1})$, initial velocity $\dot{Z}(\tau_{j-1}^+) = u_j$ and colliding elastically with two infinite-mass spherical scatterers of radius r centred at the points

$$Z(\tau_{j-1}) + r \frac{u_{j-1} - u_j}{|u_{j-1} - u_j|}, \quad \text{respectively} \quad Z(\tau_{j-2}) - r \frac{u_{j-1} - u_{j-2}}{|u_{j-1} - u_{j-2}|}.$$

These steps define the process $t \mapsto Z(t)$ in a unique way and adapted to the natural filtration of the process $t \mapsto Y(t)$. As we already stressed, the basic facts about the process $t \mapsto Z(t)$ (listed earlier in this subsection) will be proved in later sections of the paper, and will have a key role in proving our main Theorem 2.

Consistently with the notations adopted for the processes $Y(t)$ and $X(t)$, we denote $Z_k := Z(\tau_k)$ for $k \geq 0$.

3. No Mismatches up to $T \mathcal{D} o(r^{-1})$: Proof of Theorem 1

In this section we prove that the Markovian flight trajectory $Y(t)$, up to time scales of order $T = T(r) = o(r^{-1})$, is mechanically r -consistent with probability $1 - o(1)$, and therefore the coupling bound of Theorem 1 holds. On the way we establish various bounds to be used in later sections. This section is purely classical-probabilistic. It also prepares the ideas (and notation) for Sect. 5 where a similar argument is explored in more complex form.

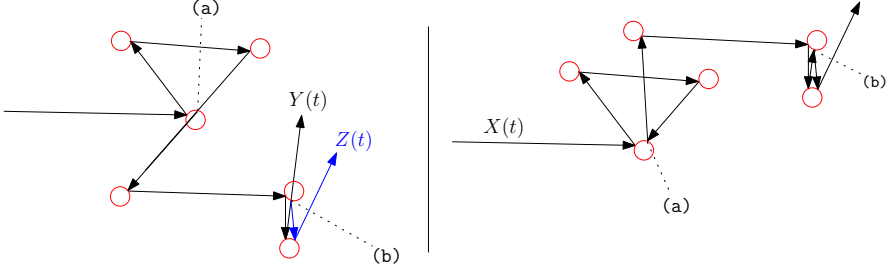


Fig. 2. The above image shows a section of trajectory during which X , Y , and Z would all three differ. On the left we see Y and Z remain together until point (b), where a direct recollision is respected by Z . Note that Z ignores the mismatch at (a) as it is indirect. On the right, the process X is coupled to Y on the left. Note that X respects the indirect recollision at point (a) and the direct recollision at (b)

3.1. Interferences. Let $t \rightarrow Y(t)$ and $t \rightarrow Y^*(t)$ be two independent Markovian flight processes. Think about $Y(t)$ as running forward and $Y^*(t)$ as running backwards in time. (Note, that the Markovian flight process has invariant law under time reversal.) Define the following events

$$\begin{aligned}\widehat{W}_j &:= \{\min\{|Y(t) - Y'_j| : 0 < t < \tau_{j-1}\} < r\}, \\ \widetilde{W}_j &:= \{\min\{|Y'_k - Y(t)| : 0 \leq k < j-1, \tau_{j-1} < t < \tau_j\} < r\}, \\ \widehat{W}_j^* &:= \{\min\{|Y^*(t) - Y'_1| : 0 < t < \tau_{j-1}\} < r\}, \\ \widetilde{W}_j^* &:= \{\min\{|Y_k^{*'} - Y(t)| : 0 < k \leq j-1, 0 < t < \tau_1\} < r\}, \\ \widehat{W}_\infty^* &:= \{\min\{|Y^*(t) - Y'_1| : 0 < t < \infty\} < r\}, \\ \widetilde{W}_\infty^* &:= \{\min\{|Y_k^{*'} - Y(t)| : 0 < k < \infty, 0 < t < \tau_1\} < r\}.\end{aligned}$$

In words \widehat{W}_j is the event that the virtual collision at Y_j is *shadowed* by the past path. While \widetilde{W}_j is the event that in the time interval (τ_{j-1}, τ_j) there is a *virtual recollision* with a past scatterer.

It is obvious that

$$\begin{aligned}\mathbf{P}(\widehat{W}_j) &= \mathbf{P}(\widehat{W}_j^*) \leq \mathbf{P}(\widehat{W}_{j+1}^*) \leq \mathbf{P}(\widehat{W}_\infty^*), \\ \mathbf{P}(\widetilde{W}_j) &= \mathbf{P}(\widetilde{W}_j^*) \leq \mathbf{P}(\widetilde{W}_{j+1}^*) \leq \mathbf{P}(\widetilde{W}_\infty^*).\end{aligned}\tag{27}$$

On the other hand, by union bound and independence

$$\begin{aligned}\mathbf{P}(\widehat{W}_\infty^*) &\leq \sum_{z \in \mathbb{Z}^3} \mathbf{P}(\{1 \leq k < \infty : Y_k^* \in B_{zr, 2r}\} \neq \emptyset) \mathbf{P}(\{0 < t \leq \xi_1 : Y(t) \in B_{zr, 2r}\} \neq \emptyset), \\ &\leq \sum_{z \in \mathbb{Z}^3} (2r)^{-1} \mathbf{E}(|\{1 < k < \infty : Y_k^* \in B_{zr, 2r}\}|) \mathbf{E}(|\{0 < t \leq \xi_1 : Y(t) \in B_{zr, 3r}\}|), \\ \mathbf{P}(\widetilde{W}_\infty^*) &\leq \sum_{z \in \mathbb{Z}^3} \mathbf{P}(\{0 < t < \infty : Y^*(t) \in B_{zr, 2r}\} \neq \emptyset) \mathbf{P}(Y_1 \in B_{zr, 2r}), \\ &\leq \sum_{z \in \mathbb{Z}^3} (2r)^{-1} \mathbf{E}(|\{0 < t < \infty : Y^*(t) \in B_{zr, 3r}\}|) \mathbf{P}(Y_1 \in B_{zr, 2r}),\end{aligned}\tag{28}$$

recall that ξ_1 is the time of the first collision for the forwards process $t \mapsto Y(t)$. Here and in the rest of the paper we use the notation $|\{\dots\}|$ for either cardinality or Lebesgue measure of the set $\{\dots\}$, depending on context.

3.2. *Occupation measures (Green's functions).* Define the following occupation measures (Green's functions): for $A \subset \mathbb{R}^3$

$$\begin{aligned} g(A) &:= \mathbf{P}(Y_1 \in A) \\ h(A) &:= \mathbf{E}(|\{0 < t \leq \xi_1 : Y(t) \in A\}|) \\ G(A) &:= \mathbf{E}(|\{1 \leq k < \infty : Y_k \in A\}|) \\ H(A) &:= \mathbf{E}(|\{0 < t < \infty : Y(t) \in A\}|). \end{aligned}$$

Obviously,

$$\begin{aligned} G(A) &= g(A) + \int_{\mathbb{R}^3} g(A - x)G(dx) \\ H(A) &= h(A) + \int_{\mathbb{R}^3} h(A - x)G(dx). \end{aligned} \tag{29}$$

3.3. *Bounds.*

Lemma 2. *The following identities and upper bounds hold:*

$$h(dx) = g(dx) \leq L(dx) \tag{30}$$

$$H(dx) = G(dx) \leq K(dx) + L(dx), \tag{31}$$

where

$$K(dx) := C \min\{1, |x|^{-1}\}dx, \quad L(dx) := Ce^{-c|x|} |x|^{-2} dx, \tag{32}$$

with appropriately chosen $C < \infty$ and $c > 0$.

Proof of Lemma 2. The identity $h = g$ is a direct consequence of the flight length ξ being $EXP(1)$ -distributed. The distribution g has the explicit expression

$$g(dx) = C |x|^{-2} e^{-|x|} dx$$

from which the the upper bound (30) follows.

(31) then follows from (29) and standard Green's function estimate for a random walk with step distribution g . \square

For later use we introduce the conditional versions—conditioned on the sequence $\underline{\epsilon}$ (see (20))—of the bounds (30), (31). In this order we define the conditional versions of the Green's functions, given $\epsilon \in \{0, 1\}$, respectively $\underline{\epsilon} \in \{0, 1\}^{\mathbb{N}}$:

$$\begin{aligned} g_\epsilon(A) &:= \mathbf{P}(Y_1 \in A \mid \epsilon) \\ h_\epsilon(A) &:= \mathbf{E}(|\{0 < t \leq \xi_1 : Y(t) \in A\}| \mid \epsilon) \\ G_{\underline{\epsilon}}(A) &:= \mathbf{E}(|\{1 \leq k < \infty : Y_k \in A\}| \mid \underline{\epsilon}) \\ H_{\underline{\epsilon}}(A) &:= \mathbf{E}(|\{0 < t < \infty : Y(t) \in A\}| \mid \underline{\epsilon}), \end{aligned}$$

and state the conditional version of Lemma 2:

Lemma 3. *The following upper bounds hold uniformly in $\epsilon \in \{0, 1\}^{\mathbb{N}}$:*

$$g_\epsilon(dx) \leq L(dx), \quad h_\epsilon(dx) \leq L(dx), \quad (33)$$

$$G_\epsilon(dx) \leq K(dx) + L(dx), \quad H_\epsilon(dx) \leq K(dx) + L(dx), \quad (34)$$

with $K(x)$ and $L(x)$ as in (32), with appropriately chosen constants $C < \infty$ and $c > 0$.

Proof of Lemma 3. Noting that

$$g_\epsilon(dx) \leq C |x|^{-2} e^{-|x|} dx, \quad h_\epsilon(dx) \leq C |x|^{-2} e^{-|x|} dx,$$

the proof of Lemma 3 follows very much the same lines as the proof of Lemma 2. We omit the details. \square

3.4. Computation. According to (27), (28), for every $j = 1, 2, \dots$

$$\begin{aligned} \mathbf{P}(\widehat{W}_j) &\leq \mathbf{P}(\widehat{W}_\infty^*) \leq (2r)^{-1} \sum_{z \in \mathbb{Z}^3} G(B_{zr, 2r}) h(B_{zr, 3r}), \\ \mathbf{P}(\widetilde{W}_j) &\leq \mathbf{P}(\widetilde{W}_\infty^*) \leq (2r)^{-1} \sum_{z \in \mathbb{Z}^3} H(B_{zr, 3r}) g(B_{zr, 2r}). \end{aligned}$$

Moreover, straightforward computations yield

Lemma 4. *In dimension $d = 3$ the following bounds hold, with some $C < \infty$*

$$\sum_{z \in \mathbb{Z}^3} K(B_{zr, 3r}) L(B_{zr, 2r}) \leq Cr^3, \quad \sum_{z \in \mathbb{Z}^3} L(B_{zr, 3r}) L(B_{zr, 2r}) \leq Cr^2. \quad (35)$$

Proof of Lemma 4. The bounds (35) readily follow from explicit computations. We omit the details. \square

We conclude this section with the following consequence of the above arguments and computations.

Corollary 1. *There exists a constant $C < \infty$ such that for any $j \geq 1$:*

$$\mathbf{P}(\widetilde{W}_j) \leq Cr, \quad \mathbf{P}(\widehat{W}_j) \leq Cr. \quad (36)$$

3.5. No mismatching—up to $T \sim o(r^{-1})$. Define the stopping time

$$\sigma := \min\{j > 0 : \max\{\mathbb{1}_{\widehat{W}_j}, \mathbb{1}_{\widetilde{W}_j}\} = 1\},$$

and note that by construction

$$\inf\{t > 0 : X(t) \neq Y(t)\} \geq \tau_{\sigma-1}. \quad (37)$$

Lemma 5. *Let $T = T(r)$ be such that $\lim_{r \rightarrow 0} T(r) = \infty$ and $\lim_{r \rightarrow 0} rT(r) = 0$. Then*

$$\lim_{r \rightarrow 0} \mathbf{P}(\tau_{\sigma-1} < T) = 0. \quad (38)$$

Proof of Lemma 5.

$$\mathbf{P}(\tau_{\sigma-1} < T) \leq \mathbf{P}(\sigma \leq 2T) + \mathbf{P}\left(\sum_{j=1}^{2T-1} \xi_j < T\right) \leq CrT + Ce^{-cT}, \quad (39)$$

where $C < \infty$ and $c > 0$. The first term in the middle expression of (39) is bounded by union bound and (36) of Corollary 1. In bounding the second term we use a large deviation upper bound for the sum of independent $EXP(1)$ -distributed ξ_j -s.

Finally, (38) readily follows from (39). \square

(8) follows directly from (37) and (38), and this concludes the proof of Theorem 1. \square

4. Beyond the Naïve Coupling

The forthcoming parts of the paper rely on the joint realisation (coupling) of the *three* processes $t \mapsto (Y(t), X(t), Z(t))$ as described in Sect. 2. In particular, recall the construction of the process $t \mapsto Z(t)$ from Sect. 2.4.

4.1. Breaking Z into legs. Let $\Gamma_0 := 0$, $\Theta_0 = 0$ and for $n \geq 1$

$$\begin{aligned} \Gamma_n &:= \min\{j \geq \Gamma_{n-1} + 2 : \min\{\xi_{j-1}, \xi_j, \xi_{j+1}, \xi_{j+2}\} > 1\}, & \gamma_n &:= \Gamma_n - \Gamma_{n-1}, \\ \Theta_n &:= \tau_{\Gamma_n}, & \theta_n &:= \Theta_n - \Theta_{n-1}, \end{aligned} \quad (40)$$

and denote

$$\begin{aligned} \xi_{n,j} &:= \xi_{\Gamma_{n-1}+j}, & u_{n,j} &:= u_{\Gamma_{n-1}+j}, & y_{n,j} &:= y_{\Gamma_{n-1}+j}, & 1 \leq j \leq \gamma_n, \\ Y_n(t) &:= Y(\Theta_{n-1} + t) - Y(\Theta_{n-1}), & & & & 0 \leq t \leq \theta_n, \\ Z_n(t) &:= Z(\Theta_{n-1} + t) - Z(\Theta_{n-1}), & & & & 0 \leq t \leq \theta_n. \end{aligned}$$

Then, it is straightforward that the *packs of random variables*

$$\varpi_n := (\gamma_n; (\xi_{n,j}, u_{n,j}) : 1 \leq j \leq \gamma_n), \quad n \geq 1, \quad (41)$$

are fully independent (for $n \geq 1$), and also identically distributed for $n \geq 2$. (The first pack is deficient if $\min\{\xi_0, \xi_1\} < 1$.) It is also straightforward that the *legs* of the Markovian flight process

$$(\theta_n; Y_n(t) : 0 \leq t \leq \theta_n), \quad n \geq 1,$$

are fully independent, and identically distributed for $n \geq 2$.

A *key observation* is that due to the rules of construction of the process $t \mapsto Z(t)$ exposed in Sect. 2.4, the legs

$$(\theta_n; Z_n(t) : 0 \leq t \leq \theta_n), \quad n \geq 1, \quad (42)$$

of the auxiliary process $t \mapsto Z(t)$ are also *independently constructed* from the packs (41), following the rules in Sect. 2.4. Note, that the restrictions $|y_{j-1}| < 1$ in (24) were imposed exactly in order to ensure this independence of the legs (42). Therefore we will construct now the auxiliary process $t \mapsto Z(t)$ and its time reversal $t \mapsto Z^*(t)$

from an infinite sequence of independent packs (41). In order to reduce unnecessary complications of notation from now on we assume $\min\{\xi_0, \xi_1\} > 1$.

Remark: In order to break up the auxiliary process $t \mapsto Z(t)$ into *independent legs* the choice of simpler stopping times

$$\Gamma'_n := \min\{j \geq \Gamma_{n-1} + 1 : \min\{\xi_j, \xi_{j+1}\} > 1\},$$

would work. However, we need the slightly more complicated stoppings Γ_n , given in (40), for some other reasons which will become clear towards the end of Sect. 4.2 and in the statement and proof of Lemma 6.

4.2. One leg. Let $\xi_j, u_j, j \geq 1$, be fully independent random variables with the distributions (18), conditioned to

$$\min\{\xi_1, \xi_2\} > 1,$$

and y_j as in (19). Let

$$\gamma := \min\{j \geq 2 : \min\{\xi_{j-1}, \xi_j, \xi_{j+1}, \xi_{j+2}\} > 1\} \in \{2\} \cup \{5, 6, \dots\}. \quad (43)$$

Note that γ can not assume the values $\{1, 3, 4\}$. Call

$$\varpi := (\gamma; (\xi_j, u_j) : 1 \leq j \leq \gamma) \quad (44)$$

a *pack*, and keep the notation $\tau_j := \sum_{k=1}^j \xi_k$, and $\theta := \tau_\gamma$.

The *forward leg*

$$(\theta; Z(t) : 0 \leq t \leq \theta)$$

is constructed from the pack ϖ according to the rules given in Sect. 2.4. We will also denote

$$Z_j := Z(\tau_j), \quad 0 \leq j \leq \gamma; \quad \bar{Z} := Z_\gamma = Z(\theta).$$

These are the discrete steps, respectively, the terminal position of the leg.

It is easy to see that the distributions of γ and θ are exponentially tight: there exist constants $C < \infty$ and $c > 0$ such that for any $s \in [0, \infty)$

$$\mathbf{P}(\gamma > s) \leq Ce^{-cs}, \quad \mathbf{P}(\theta > s) \leq Ce^{-cs}. \quad (45)$$

The *backwards leg*

$$(\theta; Z^*(t) : 0 \leq t \leq \theta)$$

is constructed from the pack ϖ as

$$Z^*(t, \varpi) := Z(\theta - t, \varpi^*) - \bar{Z}(\varpi^*),$$

where the backwards pack

$$\varpi^* := (\gamma; (\xi_{\gamma-j}, -u_{\gamma-j}) : 0 \leq j \leq \gamma - 1)$$

is the time reversion of the pack ϖ .

Note that since the velocities u_n are uniformly distributed the forward and backward packs, ϖ and ϖ^* , are identically distributed. However, since under time reversal, recollisions become shadowed scattering and vice-versa, the forward and backward processes $(t \mapsto Z(t) : 0 \leq t \leq \theta)$ and $(t \mapsto Z^*(t) : 0 \leq t \leq \theta)$ are not identically distributed. The backwards process $t \mapsto Z^*(t)$ could also be defined using an explicit step-by-step construction, similar (but not identical) to those in Sect. 2.4, but we will not rely on these step-wise rules and therefore omit their explicit formulation. Thus it suffices to take $(t \mapsto Z^*(t) : 0 \leq t \leq \theta)$ to be the time-reversal of $(t \mapsto Z(t) : 0 \leq t \leq \theta)$.

Consistent with the previous notation, we denote

$$Z_j^* := Z^*(\tau_j), \quad 0 \leq j \leq \gamma; \quad \bar{Z}^* := Z_\gamma^* = Z^*(\theta) = -\bar{Z}.$$

Note, that due to the construction rules of the forward and backward legs, their beginning, middle and ending parts

$$\begin{aligned} &(\tau_1; Z(t) : 0 \leq t \leq \tau_1), \\ &(\tau_{\gamma-1} - \tau_1; Z(\tau_1 + t) - Z(\tau_1) : 0 \leq t \leq \tau_{\gamma-1} - \tau_1), \\ &(\tau_\gamma - \tau_{\gamma-1}; Z(\tau_{\gamma-1} + t) - Z(\tau_{\gamma-1}) : 0 \leq t \leq \tau_\gamma - \tau_{\gamma-1}), \end{aligned} \quad (46)$$

are *independent*, and likewise for the backwards process Z^* ,

$$\begin{aligned} &(\tau_1; Z^*(t) : 0 \leq t \leq \tau_1), \\ &(\tau_{\gamma-1} - \tau_1; Z^*(\tau_1 + t) - Z^*(\tau_1) : 0 \leq t \leq \tau_{\gamma-1} - \tau_1), \\ &(\tau_\gamma - \tau_{\gamma-1}; Z^*(\tau_{\gamma-1} + t) - Z^*(\tau_{\gamma-1}) : 0 \leq t \leq \tau_\gamma - \tau_{\gamma-1}). \end{aligned} \quad (47)$$

This fact will be of crucial importance in the proof of Proposition 2, Sect. 5.2 below. This is the reason (alluded to in the remark at the end of Sect. 4.1) we chose the somewhat complicated stopping time as defined in (43).

4.3. Multi-leg concatenation. Let $\varpi_n = (\gamma_n; (\xi_{n,j}, u_{n,j}) : 1 \leq j \leq \gamma_n)$, $n \geq 1$, be a sequence of i.i.d *packs* (44), and denote θ_n , $(Z_n(t) : 0 \leq t \leq \theta_n)$, $(Z_{n,j} : 1 \leq j \leq \gamma_n)$, $(Z_n^*(t) : 0 \leq t \leq \theta_n)$, $(Z_{n,j}^* : 1 \leq j \leq \gamma_n)$, \bar{Z}_n , \bar{Z}_n^* the various objects defined in Sect. 4.2, specified for the n -th independent leg.

In order to construct the concatenated forward and backward processes $t \mapsto Z(t)$, $t \mapsto Z^*(t)$, $0 \leq t < \infty$, we first define for $n \in \mathbb{Z}_+$, respectively $t \in \mathbb{R}_+$

$$\begin{aligned} \Gamma_n &:= \sum_{k=1}^n \gamma_k, & v_n &:= \max\{m : \Gamma_m \leq n\}, & \{n\} &:= n - \Gamma_{v_n}, \\ \Theta_n &:= \sum_{k=1}^n \theta_k, & v_t &:= \max\{m : \Theta_m < t\}, & \{t\} &:= t - \Theta_{v_t}. \end{aligned}$$

The concatenated (multi-leg) forward and backward Z -processes are

$$\begin{aligned} \Xi_n &:= \sum_{k=1}^n \bar{Z}_k, & Z_n &:= \Xi_{v_n} + Z_{v_n+1, \{n\}}, & Z(t) &:= \Xi_{v_t} + Z_{v_t+1}(\{t\}), \\ \Xi_n^* &:= \sum_{k=1}^n \bar{Z}_k^*, & Z_n^* &:= \Xi_{v_n}^* + Z_{v_n+1, \{n\}}^*, & Z^*(t) &:= \Xi_{v_t}^* + Z_{v_t+1}^*(\{t\}). \end{aligned} \quad (48)$$

Note that Ξ_n and Ξ_n^* are random walks with independent steps; $t \mapsto Z(t)$, $0 \leq t < \infty$, is exactly the Z -process constructed in Sect. 2.4, with $Z_n = Z(\tau_n)$, $0 \leq n < \infty$. Similarly, $t \mapsto Z^*(t)$, $0 \leq t < \infty$, is the time reversal of the Z -process and $Z_n^* = Z^*(\tau_n)$, $0 \leq n < \infty$.

Theorem 2 will follow from Propositions 1 and 2 of the next two sections.

4.4. Mismatches within one leg. Given a pack $\varpi = (\gamma; (\xi_j, u_j) : 1 \leq j \leq \gamma)$ (44), and arbitrary incoming and outgoing velocities $u_0, u_{\gamma+1} \in S^2$ let $((Y(t), \mathcal{X}(t), Z(t)) : 0^- < t < \theta^+)$, be the triplet of Markovian flight process, Lorentz exploration process and auxiliary Z -process jointly constructed with these data, as described in Sects. 2.1, 2.2, respectively, 2.4. We use the notation \mathcal{X} to denote a mechanical Lorentz exploration constructed using the rules of Sect. 2.2 defined on a leg for times $t \in [0^-, \theta^+]$, and independently for different legs. By $0^- < t < \theta^+$ we mean that the incoming velocities at 0^- are given as $\dot{Y}(0^-) = \dot{\mathcal{X}}(0^-) = \dot{Z}(0^-) = u_0$ and the outgoing velocities at θ^+ are $\dot{Y}(\theta^+) = \dot{Z}(\theta^+) = u_{\gamma+1}$, while $\dot{\mathcal{X}}(\theta^+)$ is determined by the construction from Sect. 2.2. That is, $\dot{\mathcal{X}}(\theta^+) = u_{\gamma+1}$ if this last scattering is not shadowed by the trajectory $(\mathcal{X}(t) : 0 \leq t \leq \theta)$ and $\dot{\mathcal{X}}(\theta^+) = \dot{\mathcal{X}}(\theta^-)$ if it is shadowed.

Proposition 1. *There exists a constant $C < \infty$ such that for any $u_0, u_{\gamma+1} \in S^2$*

$$\mathbf{P}(\mathcal{X}(t) \not\equiv Z(t) : 0^- < t < \theta^+) \leq Cr^2 |\log r|^2. \quad (49)$$

The proof of this Proposition relies on controlling the geometry of mismatches, and is postponed until Sect. 6.

4.5. Inter-leg mismatches. Let $t \mapsto Z(t)$ be a forward Z -process built up as concatenation of legs, as exposed in Sect. 4.3 and define the following events

$$\begin{aligned} \widehat{W}_j &:= \{ \min\{|Z(t) - Z'_k| : 0 < t < \Theta_{j-1}, \quad \Gamma_{j-1} < k \leq \Gamma_j\} < r \}, \\ \widetilde{W}_j &:= \{ \min\{|Z'_k - Z(t)| : 0 \leq k < \Gamma_{j-1}, \quad \Theta_{j-1} < t < \Theta_j\} < r \}. \end{aligned} \quad (50)$$

In words \widehat{W}_j is the event that a collision occurring in the j -th leg is *shadowed* by the past path. While \widetilde{W}_j is the event that within the j -th leg the Z -trajectory bumps into a scatterer placed in an earlier leg. That is, $\widetilde{W}_j \cup \widehat{W}_j$ is precisely the event that the concatenated first $j - 1$ legs and the j -th leg are mechanically r -incompatible (see Sect. 2.3).

The following proposition indicates that on our time scales there are no “inter-leg mismatches”:

Proposition 2. *There exists a constant $C < \infty$ such that for all $j \geq 1$*

$$\mathbf{P}(\widetilde{W}_j) \leq Cr^2, \quad \mathbf{P}(\widehat{W}_j) \leq Cr^2. \quad (51)$$

The proof of Proposition 2 is the content of Sect. 5.

5. Proof of Proposition 2

This section is purely probabilistic and of similar spirit as Sect. 3. The notation used is also similar. However, similar is not identical. The various Green’s functions used here, although denoted g, h, G, H , as in Sect. 3, are similar in their rôle but not the same. The estimates on them are also different.

5.1. Occupation measures (Green's functions). Let now $t \mapsto Z^*(t)$, $0 \leq t < \infty$, be a backward Z^* -process and $t \mapsto Z(t)$, $0 \leq t \leq \theta$, a forward one-leg Z -process, assumed independent. In analogy with the events \widehat{W}_j and \widetilde{W}_j defined in (50) we define

$$\begin{aligned}\widehat{W}_j^* &:= \left\{ \min\{|Z^*(t) - Z'_k| : 0 < t < \Theta_{j-1}, \quad 0 < k \leq \gamma < r\} \right\}, \\ \widetilde{W}_j^* &:= \left\{ \min\{|Z_k^* - Z(t)| : 0 < k \leq \Gamma_{j-1}, \quad 0 < t < \theta < r\} \right\}, \\ \widehat{W}_\infty^* &:= \left\{ \min\{|Z^*(t) - Z'_k| : 0 < t < \infty, \quad 0 < k \leq \gamma < r\} \right\}, \\ \widetilde{W}_\infty^* &:= \left\{ \min\{|Z_k^* - Z(t)| : 0 < k < \infty, \quad 0 < t < \theta < r\} \right\}.\end{aligned}$$

It is obvious that

$$\begin{aligned}\mathbf{P}(\widehat{W}_j) &= \mathbf{P}(\widehat{W}_j^*) \leq \mathbf{P}(\widehat{W}_{j+1}^*) \leq \mathbf{P}(\widehat{W}_\infty^*), \\ \mathbf{P}(\widetilde{W}_j) &= \mathbf{P}(\widetilde{W}_j^*) \leq \mathbf{P}(\widetilde{W}_{j+1}^*) \leq \mathbf{P}(\widetilde{W}_\infty^*).\end{aligned}\tag{52}$$

On the other hand, by the union bound and independence we have

$$\begin{aligned}\mathbf{P}(\widehat{W}_\infty^*) &\leq \sum_{z \in \mathbb{Z}^3} \mathbf{P}(\{0 < t < \infty : Z^*(t) \in B_{zr, 2r}\} \neq \emptyset) \mathbf{P}(\{1 \leq k \leq \gamma : Z_k \in B_{zr, 2r}\} \neq \emptyset) \\ &\leq \sum_{z \in \mathbb{Z}^3} (2r)^{-1} \mathbf{E}(|\{0 < t < \infty : Z^*(t) \in B_{zr, 3r}\}|) \mathbf{E}(|\{1 \leq k \leq \gamma : Z_k \in B_{zr, 2r}\}|) \\ \mathbf{P}(\widetilde{W}_\infty^*) &\leq \sum_{z \in \mathbb{Z}^3} \mathbf{P}(\{1 \leq k < \infty : Z_k^* \in B_{zr, 2r}\} \neq \emptyset) \mathbf{P}(\{0 < t \leq \theta : Z(t) \in B_{zr, 2r}\} \neq \emptyset), \\ &\leq \sum_{z \in \mathbb{Z}^3} (2r)^{-1} \mathbf{E}(|\{1 \leq k < \infty : Z_k^* \in B_{zr, 2r}\}|) \mathbf{E}(|\{0 < t \leq \theta : Z(t) \in B_{zr, 3r}\}|).\end{aligned}\tag{53}$$

Therefore, in view of (52) we have to control the mean occupation time measures appearing on the right hand side of (53).

Define the following mean occupation measures (Green's functions): for $A \subset \mathbb{R}^3$ let

$$\begin{aligned}g(A) &:= \mathbf{E}(|\{1 \leq k \leq \gamma : Z_k \in A\}|), \\ g^*(A) &:= \mathbf{E}(|\{1 \leq k \leq \gamma : Z_k^* \in A\}|), \\ h(A) &:= \mathbf{E}(|\{0 < t \leq \theta : Z(t) \in A\}|), \\ h^*(A) &:= \mathbf{E}(|\{0 < t \leq \theta : Z^*(t) \in A\}|), \\ R^*(A) &:= \mathbf{E}(|\{1 \leq n < \infty : \Xi_n^* \in A\}|), \\ G^*(A) &:= \mathbf{E}(|\{1 \leq k < \infty : Z_k^* \in A\}|), \\ H^*(A) &:= \mathbf{E}(|\{0 < t < \infty : Z^*(t) \in A\}|).\end{aligned}$$

It is obvious that

$$\begin{aligned}G^*(A) &= g^*(A) + \int_{\mathbb{R}^3} g^*(A - x) R^*(dx), \\ H^*(A) &= h^*(A) + \int_{\mathbb{R}^3} h^*(A - x) R^*(dx).\end{aligned}\tag{54}$$

5.2. Bounds.

Lemma 6. *The following upper bounds hold:*

$$\max\{g(dx), g^*(dx)\} \leq M(dx), \quad \max\{h(dx), h^*(dx)\} \leq L(dx), \quad (55)$$

$$R^*(dx) \leq K(dx), \quad (56)$$

$$G^*(dx) \leq K(dx), \quad H^*(dx) \leq K(dx) + L(dx), \quad (57)$$

where

$$K(dx) := C \min\{1, |x|^{-1}\} dx, \quad L(dx) := C e^{-c|x|} |x|^{-2} dx, \quad M(dx) := C e^{-c|x|} dx,$$

with appropriately chosen $C < \infty$ and $c > 0$.

Proof of Lemma 6. The proof of the bounds (55) hinges on the decompositions (46) and (47) of the forward and backward legs into independent parts.

Let

$$\begin{aligned} g_1(A) &:= \mathbf{P}(Z_1 \in A) = \mathbf{P}(Z_1^* \in A) \\ &= C \int_A \mathbb{1}(|x| > 1) e^{-|x|} dx, \\ h_1(A) &:= \mathbf{E}(|\{t \leq \tau_1 : Z(t) \in A\}|) = \mathbf{E}(|\{t \leq \tau_1 : Z^*(t) \in A\}|) \\ &= C' \int_A |x|^{-2} e^{-\max\{1, |x|\}} dx, \end{aligned} \quad (58)$$

and

$$\begin{aligned} g_2(A) &:= \mathbf{E}(|\{1 \leq k \leq \gamma : Z_k - Z_1 \in A\}|), \\ g_2^*(A) &:= \mathbf{E}(|\{1 \leq k \leq \gamma : Z_k^* - Z_1^* \in A\}|), \\ h_2(A) &:= \mathbf{E}(|\{0 < t \leq \theta - \tau_1 : Z(\tau_1 + t) - Z_1 \in A\}|), \\ h_2^*(A) &:= \mathbf{E}(|\{0 < t \leq \theta - \tau_1 : Z^*(\tau_1 + t) - Z_1^* \in A\}|). \end{aligned}$$

Due to the exponential tail of the distribution of γ and θ , (45), there are constants $C < \infty$ and $c > 0$ such that for any $s < \infty$

$$\begin{aligned} \max\{g_2(\{x : |x| > s\}), g_2^*(\{x : |x| > s\})\} &\leq C e^{-cs}, \\ \max\{h_2(\{x : |x| > s\}), h_2^*(\{x : |x| > s\})\} &\leq C e^{-cs}, \end{aligned} \quad (59)$$

and furthermore,

$$\begin{aligned} g_2(\mathbb{R}^3) &= g_2^*(\mathbb{R}^3) = \mathbf{E}(\gamma) < \infty, \\ h_2(\mathbb{R}^3) &= h_2^*(\mathbb{R}^3) = \mathbf{E}(\theta - \tau_1) < \infty. \end{aligned} \quad (60)$$

From the independent decompositions (47) and (46) it follows that

$$\begin{aligned} g(A) &= \int_{\mathbb{R}^3} g_2(A - x) g_1(dx), & g^*(A) &= \int_{\mathbb{R}^3} g_2^*(A - x) g_1(dx), \\ h(A) &= \int_{\mathbb{R}^3} h_2(A - x) g_1(dx) + h_1(A), & h^*(A) &= \int_{\mathbb{R}^3} h_2^*(A - x) g_1(dx) + h_1(A). \end{aligned} \quad (61)$$

The bounds (55) readily follow from the explicit expressions (58), the convolutions (61) and the bounds (59) and (60).

The bound (56) is a straightforward Green's function bound for the the random walk Ξ_n^* defined in (48), by noting that the distribution of the i.i.d. steps \bar{Z}_k^* of this random walk has bounded density and exponential tail decay.

Finally, the bounds (57) follow from the convolutions (54) and the bounds (55), (56). \square

Remark: On the difference between Lemmas 2 and 6 . Note the difference between the upper bounds for g in (30), respectively, (55), and on G in (31), respectively, (57). These are important and are due to the fact that the length first step in a Z - or Z^* -leg is distributed as $(\xi \mid \xi > 1) \sim EXP(1|0)$ rather than $\xi \sim EXP(1)$.

5.3. *Computation.* According to (53)

$$\begin{aligned} \mathbf{P}(\tilde{W}_j) &\leq \mathbf{P}(\tilde{W}_\infty^*) \leq (2r)^{-1} \sum_{z \in \mathbb{Z}^3} H^*(B_{zr,3r}) g(B_{zr,2r}), \\ \mathbf{P}(\hat{W}_j) &\leq \mathbf{P}(\hat{W}_\infty^*) \leq (2r)^{-1} \sum_{z \in \mathbb{Z}^3} G^*(B_{zr,2r}) h(B_{zr,3r}). \end{aligned} \quad (62)$$

Lemma 7. *In dimension $d = 3$ the following bounds hold, with some $C < \infty$*

$$\sum_{z \in \mathbb{Z}^3} K(B_{zr,3r}) M(B_{zr,2r}) \leq Cr^3, \quad \sum_{z \in \mathbb{Z}^3} M(B_{zr,3r}) L(B_{zr,2r}) \leq Cr^3. \quad (63)$$

Proof of Lemma 7. The bounds (63) (similarly to the bounds (35)) readily follow from explicit computations which we omit. \square

Proof of Proposition 2. Proposition 2 now follows by inserting the bounds (63) and one of the bounds in (35) into equations (62). \square

6. Proof of Proposition 1

Given a pack $\varpi = (\gamma; (\xi_j, u_j) : 1 \leq j \leq \gamma)$ (44), and arbitrary $u_0, u_{\gamma+1} \in S^2$, let $((Y(t), \mathcal{X}(t), Z(t)) : 0 \leq t \leq \theta)$ be the triplet of Markovian flight process, Lorentz exploration process and auxiliary Z -process jointly constructed with these data. We will prove the following bounds, stated in increasing order of difficulty/complexity.

$$\mathbf{P}(\{\mathcal{X}(t) \neq Z(t) : 0^- \leq t \leq \theta^+\} \cap \{\sum_{j=1}^{\gamma} \eta_j > 1\}) \leq Cr^2 |\log r|, \quad (64)$$

$$\mathbf{P}(\{\mathcal{X}(t) \neq Z(t) : 0^- \leq t \leq \theta^+\} \cap \{\sum_{j=1}^{\gamma} \eta_j = 0\}) \leq Cr^2 |\log r|, \quad (65)$$

$$\mathbf{P}(\{\mathcal{X}(t) \neq Z(t) : 0^- \leq t \leq \theta^+\} \cap \{\sum_{j=1}^{\gamma} \eta_j = 1\}) \leq Cr^2 |\log r|^2. \quad (66)$$

Note that by construction $\eta_1 = \eta_2 = \eta_3 = \eta_\gamma = 0$, so the sums on the left hand side go actually from 4 to $\gamma - 1$. We stated and prove these bounds in their increasing order of complexity: (64) (proved in Sect. 6.1) and (65) (proved in Sect. 6.2) are of purely probabilistic nature while (66) (proved in Sects. 6.3–6.7) also relies on the finer geometric understanding of the mismatch events $\hat{\eta}_j = 1$ and $\tilde{\eta}_j = 1$.

6.1. Proof of (64). This follows directly from Lemma 1. Indeed, given γ and $\underline{\epsilon} = (\epsilon_j)_{1 \leq j \leq \gamma}$, due to (26),

$$\begin{aligned} \mathbf{P}\left(\sum_{j=1}^{\gamma} \eta_j > 1 \mid \underline{\epsilon}\right) &\leq \gamma \max_j \mathbf{P}(\eta_j = \eta_{j+1} = 1 \mid \underline{\epsilon}) + \frac{\gamma^2}{2} \max_{j,k: |j-k| > 1} \mathbf{P}(\eta_j = \eta_k = 1 \mid \underline{\epsilon}) \\ &\leq C\gamma r^2 |\log r| + C\gamma^2 r^2, \end{aligned}$$

and hence, due to the exponential tail bound (45) we get

$$\mathbf{P}\left(\sum_{j=4}^{\gamma-1} \eta_j > 1\right) = \mathbf{E}\left(\mathbf{P}\left(\sum_{j=4}^{\gamma-1} \eta_j > 1 \mid \underline{\epsilon}\right)\right) \leq Cr^2 |\log r|,$$

which concludes the proof of (64). \square

6.2. Proof of (65). First note that by construction of the processes $((\mathcal{X}(t), Z(t)) : 0^- < t < \theta^+)$ the following identities hold:

$$\begin{aligned} \{\mathcal{X}(t) \neq Z(t) : 0^- \leq t \leq \theta^+\} \cap \left\{\sum_{j=1}^{\gamma} \eta_j = 0\right\} &= \{\mathcal{X}(t) \neq Y(t) : 0^- \leq t \leq \theta^+\} \cap \left\{\sum_{j=1}^{\gamma} \eta_j = 0\right\}, \\ \{\mathcal{X}(t) \neq Y(t) : 0^- \leq t \leq \theta^+\} &= \bigcup_{0 < j < \gamma} \left\{ \min_{\tau_j \leq t \leq \theta} |Y'_{j-1} - Y(t)| < r \right\} \cup \left\{ \min_{0 \leq t \leq \tau_j} |Y'_{j+1} - Y(t)| < r \right\}, \end{aligned}$$

and, hence

$$\begin{aligned} &\{\mathcal{X}(t) \neq Z(t) : 0^- \leq t \leq \theta^+\} \cap \left\{\sum_{j=1}^{\gamma} \eta_j = 0\right\} \\ &= \bigcup_{0 < j < \gamma} \left(\left\{ \min_{\tau_j \leq t \leq \tau_{j+1}} |Y'_{j-1} - Y(t)| < r \right\} \cup \left\{ \min_{\tau_{j-1} \leq t \leq \tau_j} |Y'_{j+1} - Y(t)| < r \right\} \right) \cap \{\xi_j > 1\} \\ &\cup \bigcup_{0 < j < \gamma} \left(\left\{ \min_{\tau_{j+1} \leq t \leq \theta} |Y'_{j-1} - Y(t)| < r \right\} \cup \left\{ \min_{0 \leq t \leq \tau_{j-1}} |Y'_{j+1} - Y(t)| < r \right\} \right), \\ &\subset \bigcup_{0 < j < \gamma} \left(\left\{ \min_{\tau_j \leq t \leq \tau_{j+1}} |Y_{j-1} - Y(t)| < 2r \right\} \cup \left\{ \min_{\tau_{j-1} \leq t \leq \tau_j} |Y_{j+1} - Y(t)| < 2r \right\} \right) \cap \{\xi_j > 1\} \\ &\cup \bigcup_{0 < j < \gamma} \left(\left\{ \min_{\tau_{j+1} \leq t \leq \theta} |Y_{j-1} - Y(t)| < 2r \right\} \cup \left\{ \min_{0 \leq t \leq \tau_{j-1}} |Y_{j+1} - Y(t)| < 2r \right\} \right). \end{aligned} \quad (67)$$

By simple geometric inspection we see

$$\left\{ \min_{\tau_j \leq t \leq \tau_{j+1}} |Y_{j-1} - Y(t)| < 2r \right\} \cap \{\xi_j > 1\} \subset \{\angle(-u_{j-1}, u_j) < 4r\},$$

$$\left\{ \min_{\tau_{j-1} \leq t \leq \tau_j} |Y_{j+1} - Y(t)| < 2r \right\} \cap \{\xi_j > 1\} \subset \{\angle(-u_{j+1}, u_j) < 4r\}.$$

And therefore,

$$\begin{aligned} \max_{\underline{\epsilon}} \mathbf{P}\left(\left\{ \min_{\tau_j \leq t \leq \tau_{j+1}} |Y_{j-1} - Y(t)| < 2r \right\} \cap \{\xi_j > 1\} \mid \underline{\epsilon}\right) &\leq Cr^2 \\ \max_{\underline{\epsilon}} \mathbf{P}\left(\left\{ \min_{\tau_{j-1} \leq t \leq \tau_j} |Y_{j+1} - Y(t)| < 2r \right\} \cap \{\xi_j > 1\} \mid \underline{\epsilon}\right) &\leq Cr^2. \end{aligned} \quad (68)$$

On the other hand, from the conditional Green's function computations of Sect. 3, in particular from Lemma 3, we get

$$\begin{aligned} \max_{\underline{\epsilon}} \mathbf{P}\left(\min_{\tau_{j+1} \leq t \leq \theta} |Y_{j-1} - Y(t)| < 2r \mid \underline{\epsilon}\right) &\leq \sup_{\underline{\epsilon}} \mathbf{P}\left(\min_{\tau_2 \leq t < \infty} |Y(t)| < 2r \mid \underline{\epsilon}\right) \leq Cr^2 |\log r|, \\ \max_{\underline{\epsilon}} \mathbf{P}\left(\min_{0 \leq t \leq \tau_{j-1}} |Y_{j+1} - Y(t)| < 2r \mid \underline{\epsilon}\right) &\leq \sup_{\underline{\epsilon}} \mathbf{P}\left(\min_{\tau_2 \leq t < \infty} |Y(t)| < 2r \mid \underline{\epsilon}\right) \leq Cr^2 |\log r|. \end{aligned} \quad (69)$$

Putting (67), (68) and (69) together yields

$$\mathbf{P}(\{\mathcal{X}(t) \neq Z(t) : 0^- \leq t \leq \theta^+\} \cap \{\sum_{j=4}^{\gamma-1} \eta_j = 0\} \mid \underline{\epsilon}) \leq C\gamma r^2 |\log r|,$$

and hence, taking expectation over $\underline{\epsilon}$, we get (65).

6.3. Proof of (66)—preparations. Let $\gamma \in \{2\} \cup \{5, 6, \dots\}$, and $\underline{\epsilon} = (\epsilon_j)_{1 \leq j \leq \gamma} \in \{0, 1\}^\gamma$ compatible with the definition of a pack, and $3 < k < \gamma$ be fixed. Given a pack ϖ with signature $\underline{\epsilon}$ we define yet another auxiliary process $(Z^{(k)}(t) : 0^- < t < \theta^+)$ as follows:

- On $0^- < t \leq \tau_{k-1}$, $Z^{(k)}(t) = Y(t)$.
- On $\tau_{k-1} < t \leq \tau_k$, $Z^{(k)}(t)$ is constructed according to the rules of the Z -process, given in Sect. 2.4.
- On $\tau_k < t < \theta^+$, $Z^{(k)}(t) = Z^{(k)}(\tau_k) + Y(t) - Y(\tau_k)$.

Note that on the event $\{\eta_j = \delta_{j,k} : 1 \leq j \leq \gamma\}$ we have $Z^{(k)}(t) \equiv Z(t)$, $0^- < t < \theta^+$.

We will show that

$$\begin{aligned} \max_{\underline{\epsilon}, k} \mathbf{P}(\{\mathcal{X}(t) \neq Z^{(k)}(t) : 0^- \leq t \leq \theta^+\} \cap \{\eta_j = \delta_{j,k} : 1 \leq j \leq \gamma\} \mid \underline{\epsilon}) \\ \leq \max_{\underline{\epsilon}, k} \mathbf{P}(\{\mathcal{X}(t) \neq Z^{(k)}(t) : 0^- \leq t \leq \theta^+\} \cap \{\eta_k = 1\} \mid \underline{\epsilon}) \\ \leq C\gamma^2 r^2 |\log r|^2, \end{aligned} \quad (70)$$

and hence

$$\begin{aligned} \max_{\underline{\epsilon}} \mathbf{P}(\{\mathcal{X}(t) \neq Z(t) : 0^- \leq t \leq \theta^+\} \cap \{\sum_{k=1}^{\gamma} \eta_k = 1\} \mid \underline{\epsilon}) \\ \leq \gamma \max_{\underline{\epsilon}, k} \mathbf{P}(\{\mathcal{X}(t) \neq Z(t) : 0^- \leq t \leq \theta^+\} \cap \{\eta_j = \delta_{j,k} : 1 \leq j \leq \gamma\} \mid \underline{\epsilon}) \\ \leq C\gamma^3 r^2 |\log r|^2. \end{aligned}$$

Then, taking expectation over $\underline{\epsilon}$ we get (66).

In order to prove (70) first write

$$\begin{aligned} & \mathbf{P}(\{\mathcal{X}(t) \neq Z^{(k)}(t) : 0^- \leq t \leq \theta^+\} \cap \{\eta_j = \delta_{j,k} : 1 \leq j \leq \gamma\} \mid \underline{\epsilon}) \\ & \leq \mathbf{P}(\{\mathcal{X}(t) \neq Z^{(k)}(t) : 0^- \leq t \leq \theta^+\} \cap \{\eta_k = 1\} \mid \underline{\epsilon}) \\ & = \mathbf{P}(\{\mathcal{X}(t) \neq Z^{(k)}(t) : 0^- \leq t \leq \theta^+\} \cap \{\widehat{\eta}_k = 1\} \mid \underline{\epsilon}) \\ & + \mathbf{P}(\{\mathcal{X}(t) \neq Z^{(k)}(t) : 0^- \leq t \leq \theta^+\} \cap \{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\} \mid \underline{\epsilon}), \end{aligned}$$

and note that the three parts

$$\begin{aligned} (Z^{(k)}(t) : 0^- < t < \tau_{k-3}) &= (Y(t) : 0^- < t < \tau_{k-3}), \\ (Z^{(k)}(\tau_{k-3} + t) - Z^{(k)}(\tau_{k-3}) : 0 \leq t \leq \tau_k - \tau_{k-3}), \\ (Z^{(k)}(\tau_k + t) - Z^{(k)}(\tau_k) : 0 \leq t < \theta^+ - \tau_k) &= (Y(\tau_k + t) - Y(\tau_k) : 0 \leq t < \theta^+ - \tau_k), \end{aligned} \quad (71)$$

are *independent*—even if the events $\{\widehat{\eta}_k = 1\}$, respectively, $\{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\}$ are specified.

From the construction of the processes $((\mathcal{X}(t), Z^{(k)}(t)) : 0^- < t < \theta^+)$ it follows that if $(Z^{(k)}(t) : 0^- < t < \theta^+)$ is mechanically r -consistent then $(\mathcal{X}(t) \equiv Z^{(k)}(t) : 0^- < t < \theta^+)$.

Denote by $A_{a,a}^{(k)}$, $1 \leq a \leq 3$, the event that the a -th part of the decomposition (71) is mechanically r -inconsistent, and by $A_{a,b} = A_{b,a}$, $1 \leq a, b \leq 3$, $a \neq b$, the event that the a -th and b -th parts of the decomposition (71) are mechanically r -incompatible—in the sense of the definitions (22) and (23) in Sect. 2.3. In order to prove (70) we will have to prove appropriate upper bounds on the conditional probabilities

$$\begin{aligned} & \mathbf{P}(\{\widehat{\eta}_k = 1\} \cap A_{a,b}^{(k)} \mid \underline{\epsilon}), \\ & \mathbf{P}(\{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\} \cap A_{a,b}^{(k)} \mid \underline{\epsilon}), \end{aligned} \quad a, b = 1, 2, 3. \quad (72)$$

These are altogether 12 bounds. However, some of them are formally very similar.

$A_{1,1}^{(k)}$, $A_{3,3}^{(k)}$ and $A_{1,3}^{(k)}$ do not involve the middle part and therefore do not rely on the geometric arguments of the forthcoming Sects. 6.4–6.6. Applying directly (25), (33), (35) and similar procedures as in Sect. 3.4, without any new effort we get

$$\begin{aligned} & \mathbf{P}(\{\widehat{\eta}_k = 1\} \cap A_{a,b}^{(k)} \mid \underline{\epsilon}) \leq C\gamma^2 r^2, \\ & \mathbf{P}(\{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\} \cap A_{a,b}^{(k)} \mid \underline{\epsilon}) \leq C\gamma^2 r^2, \end{aligned} \quad a, b = 1, 3. \quad (73)$$

We omit the repetition of these details.

The remaining six bounds rely on the geometric arguments of Sects. 6.4–6.6 and, therefore, are postponed to Sect. 6.7.

6.4. Geometric estimates. We analyse the middle segment of the process $Z^{(k)}$, presented in (71), restricted to the events $\{\widehat{\eta}_k = 1\}$, respectively, $\{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\}$. Since everything done in this analysis is invariant under time and space translations and also under rigid rotations of \mathbb{R}^3 it will be notationally convenient to place the origin of space-time at $(\tau_{k-2}, Z(\tau_{k-2}))$ and choose $u_{k-2} = e = (1, 0, 0)$, a fixed element of S^2 . So,

the ingredient random variables are $(\xi_-, u, \xi, v, \xi_+)$, fully independent and distributed as $\xi_- \sim EXP(1|\epsilon_{k-2})$, $\xi \sim EXP(1|\epsilon_{k-1}) = EXP(1|1)$, $\xi_+ \sim EXP(1|\epsilon_k)$, $u, v \sim UNI(S^2)$.

It will be enlightening to group the ingredient variables as $(\xi_-, (u, \xi, v), \xi_+)$, and accordingly write the sample space of this reduced context as $\mathbb{R}_+ \times \mathbb{D} \times \mathbb{R}_+$, where $\mathbb{D} := S^2 \times \mathbb{R}_+ \times S^2$, with the probability measure $EXP(1|\epsilon_{k-2}) \times \mu \times EXP(1|\epsilon_k)$ where, on \mathbb{D} ,

$$\mu = UNI(S^2) \times EXP(1|1) \times UNI(S^2). \quad (74)$$

For $r < 1$, let $\widehat{\sigma}_r, \widetilde{\sigma}_r : \mathbb{D} \rightarrow \mathbb{R}_+ \cup \{\infty\}$ be

$$\begin{aligned} \widehat{\sigma}_r(u, \xi, v) &:= \inf\{t : \left| \xi u + r \frac{u-v}{|u-v|} + te \right| < r\}, \\ \widetilde{\sigma}_r(u, \xi, v) &:= \inf\{t : \left| \xi u + r \frac{u-e}{|u-e|} + tv \right| < r\}, \end{aligned}$$

(with the usual convention $\inf \emptyset = \infty$), and

$$\widehat{\mathbb{A}}_r := \{(u, \xi, v) \in \mathbb{D} : \widehat{\sigma}_r < \infty\}, \quad \widetilde{\mathbb{A}}_r := \{(u, \xi, v) \in \mathbb{D} : \widetilde{\sigma}_r < \infty\}.$$

We define the process $(\widehat{Z}_r(t) : -\infty < t < \infty)$ and $(\widetilde{Z}_r(t) : -\infty < t < \infty)$ in terms of $(u, \xi, v) \in \widehat{\mathbb{A}}_r$, respectively, $(u, \xi, v) \in \widetilde{\mathbb{A}}_r$ as follows. Strictly speaking, these are *deficient* processes, since $\mu(\widehat{\mathbb{A}}_r) < 1$, and $\mu(\widetilde{\mathbb{A}}_r) < 1$.

- On $-\infty < t \leq 0$, $\widehat{Z}_r(t) = \widetilde{Z}_r(t) = te$.
- On $0 \leq t \leq \xi$, $\widehat{Z}_r(t) = \widetilde{Z}_r(t) = tu$,
- On $\xi \leq t < \infty$,
 - $\widehat{Z}_r(t) = \widehat{Z}_r(\xi) + (t - \xi)u$,
 - $\widetilde{Z}_r(t)$ is the trajectory of a mechanical particle, with initial position $\widetilde{Z}_r(\xi)$ and initial velocity $\dot{\widetilde{Z}}_r(\xi^+) = v$, bouncing elastically between two infinite-mass spherical scatterers centred at $r \frac{e-u}{|e-u|}$, respectively, $\xi u + r \frac{u-v}{|u-v|}$, and, eventually, flying indefinitely with constant terminal velocity.

The *trapping time* $\widehat{\beta}_r, \widetilde{\beta}_r \in \mathbb{R}_+$ and *escape (terminal) velocity* $\widehat{w}_r, \widetilde{w}_r \in S^2$ of the process $\widehat{Z}_r(t)$, respectively, $\widetilde{Z}_r(t)$, are

$$\begin{aligned} \widehat{\beta}_r &:= 0, & \widehat{w}_r &:= u, \\ \widetilde{\beta}_r &:= \sup\{s < \infty : \dot{\widetilde{Z}}_r(\xi + s^+) \neq \dot{\widetilde{Z}}_r(\xi + s^-)\}, & \widetilde{w}_r &:= \dot{\widetilde{Z}}_r(\xi + \widetilde{\beta}_r^+). \end{aligned} \quad (75)$$

Note that $\widetilde{\beta}_r \geq \widetilde{\sigma}_r$.

The relation of the middle segment of (71) to \widehat{Z}_r and \widetilde{Z}_r is the following:

$$\begin{aligned} & \left(\{\widehat{\eta}_k = 1\}, (Z^{(k)}(\tau_{k-2} + t) - Z^{(k)}(\tau_{k-2}) : -\xi_{k-2} \leq t \leq \xi_{k-1} + \xi_k) \right) \\ & \sim \left(\{\xi_- > \widehat{\sigma}_r\}, (\widehat{Z}_r(t) : -\xi_- \leq t \leq \xi + \xi_+) \right), \\ & \left(\{\widehat{\eta}_k = 0\} \cap \{\widetilde{\eta}_k = 1\}, (Z^{(k)}(\tau_{k-2} + t) - Z^{(k)}(\tau_{k-2}) : -\xi_{k-2} \leq t \leq \xi_{k-1} + \xi_k) \right) \\ & \sim \left(\{\xi_- \leq \widehat{\sigma}_r\} \cap \{\xi_+ > \widetilde{\sigma}_r\}, (\widetilde{Z}_r(t) : -\xi_- \leq t \leq \xi + \xi_+) \right), \end{aligned} \quad (76)$$

where \sim stands for equality in distribution. (Note, that the sequence signatures $(\epsilon_n)_{n \geq 1}$ is determined by the sequence of flight times $(\xi_n)_{n \geq 1}$.) So, in order to prove (70) we have to prove some subtle estimates for the processes \tilde{Z}_r and $\tilde{\tilde{Z}}_r$. The main estimates are collected in Proposition 3 below.

Proposition 3. *There exists a constant $C < \infty$, such that for all $r < 1$ and $s \in (0, \infty)$, the following bounds hold:*

$$\mu((u, h, v) \in \hat{\mathbb{A}}_r : \angle(-e, \hat{w}_r) < s) \leq Cr \min\{s, 1\}, \quad (77)$$

$$\mu((u, h, v) \in \tilde{\mathbb{A}}_r : \angle(-e, \tilde{w}_r) < s) \leq Cr \min\{s(|\log s| \vee 1), 1\}, \quad (78)$$

$$\mu((u, h, v) \in \tilde{\tilde{\mathbb{A}}}_r : r^{-1}\tilde{\beta}_r > s) \leq Cr \min\{s^{-1}(|\log s| \vee 1), 1\}. \quad (79)$$

Remarks: The bound (77) is sharp in the sense that a lower bound of the same order can be proved. In contrast, we think that the upper bound in (78) is not quite sharp. However, it is sufficient for our purposes so we don't strive for a better estimate.

The following consequence of Proposition 3 will be used to prove (66).

Corollary 2. *There exists a constant $C < \infty$ such that the following bounds hold:*

$$\mathbf{P}(\{\hat{\eta}_k = 1\} \cap \{\min_{\tau_{k-2} \leq t \leq \tau_k} |Z^{(k)}(t) - Z^{(k)}(\tau_{k-3})| < s\} \mid \epsilon) \leq Crs(|\log s| \vee 1), \quad (80)$$

$$\mathbf{P}(\{\hat{\eta}_k = 1\} \cap \{\min_{\tau_{k-3} \leq t \leq \tau_{k-1}} |Z^{(k)}(t) - Z^{(k)}(\tau_k)| < s\} \mid \epsilon) \leq Crs(|\log s| \vee 1), \quad (81)$$

$$\begin{aligned} & \mathbf{P}(\{\hat{\eta}_k = 0\} \cap \{\tilde{\eta}_k = 1\} \cap \{\min_{\tau_{k-2} \leq t \leq \tau_k} |Z^{(k)}(t) - Z^{(k)}(\tau_{k-3})| < s\} \mid \epsilon) \\ & \leq Cr \max\{s |\log s|^2, r |\log r|^2\}, \end{aligned} \quad (82)$$

$$\begin{aligned} & \mathbf{P}(\{\hat{\eta}_k = 0\} \cap \{\tilde{\eta}_k = 1\} \cap \{\min_{\tau_{k-3} \leq t \leq \tau_{k-1} + \tilde{\beta}} |Z^{(k)}(t) - Z^{(k)}(\tau_k)| < s\} \mid \epsilon) \\ & \leq Cr \max\{s |\log s|^2, r |\log r|^2\}. \end{aligned} \quad (83)$$

Proposition 3 and its Corollary 2 are proved in Sect. 6.5, respectively, 6.6.

6.5. Geometric estimates ctd: Proof of Proposition 3.

6.5.1. Preparations Beside the probability measure μ (see (74)) we will also need the flat Lebesgue measure on \mathbb{D} ,

$$\lambda = \text{UNI}(S^2) \times \text{LEB}(\mathbb{R}_+) \times \text{UNI}(S^2),$$

so that

$$d\mu(u, h, v) = \frac{e^{1-h}}{e-1} \mathbb{1}\{0 \leq h < 1\} d\lambda(u, h, v).$$

For $r > 0$ we define the *dilation map* $D_r : \mathbb{D} \rightarrow \mathbb{D}$ as

$$D_r(u, h, v) = (u, rh, v),$$

and note that

$$\widehat{\mathbb{A}}_r = D_r \widehat{\mathbb{A}}_1 \quad \widetilde{\mathbb{A}}_r = D_r \widetilde{\mathbb{A}}_1.$$

In the forthcoming steps all events in $\widehat{\mathbb{A}}_r$ and $\widetilde{\mathbb{A}}_r$ will be mapped by the inverse dilation $D_r^{-1} = D_{r^{-1}}$ into $\widehat{\mathbb{A}}_1$, respectively, $\widetilde{\mathbb{A}}_1$. Therefore, in order to simplify notation we will use $\widehat{\mathbb{A}} := \widehat{\mathbb{A}}_1$ and $\widetilde{\mathbb{A}} := \widetilde{\mathbb{A}}_1$.

The dilation D_r transforms the measures μ as follows. Given an event $E \subset \mathbb{D}$,

$$\begin{aligned} \mu(D_r E) &= \int_{D_r E} \frac{e^{1-h}}{e-1} \mathbb{1}_{\{0 \leq h \leq 1\}} d\lambda(u, h, v) \\ &= r \int_E \frac{e^{1-rh}}{e-1} \mathbb{1}_{\{0 \leq h \leq r^{-1}\}} d\lambda(u, h, v), \end{aligned} \quad (84)$$

and hence, for any event $E \subset \mathbb{D}$ and any $\bar{h} < \infty$

$$\frac{e^{1-r\bar{h}}}{e-1} r \lambda(E \cap \{h \leq \bar{h}\}) \leq \mu(D_r E) \leq \frac{e}{e-1} r \lambda(E). \quad (85)$$

The following simple observation is of paramount importance in the forthcoming arguments:

Proposition 4. *In dimension 3 (and more)*

$$\lambda(\widehat{\mathbb{A}}) = \lambda(\widetilde{\mathbb{A}}) < \infty. \quad (86)$$

Proof of Proposition 4. Obviously,

$$\begin{aligned} \widehat{\mathbb{A}} \subset \widehat{\mathbb{A}}' &:= \{(u, h, v) \in \mathbb{D} : \angle(-e, u) \leq 2h^{-1}\}, \\ \widetilde{\mathbb{A}} \subset \widetilde{\mathbb{A}}' &:= \{(u, h, v) \in \mathbb{D} : \angle(-u, v) \leq 2h^{-1}\}. \end{aligned}$$

Since, in dimension 3,

$$\begin{aligned} \left| \{(u, v) \in S^2 \times S^2 : \angle(-e, u) < 2h^{-1}\} \right| &= \\ \left| \{(u, v) \in S^2 \times S^2 : \angle(-u, v) < 2h^{-1}\} \right| &\leq C \min\{h^{-2}, 1\}, \end{aligned}$$

the claim follows by integrating over $h \in \mathbb{R}_+$. \square

Remark: In 2-dimension, the corresponding sets $\widehat{\mathbb{A}}$, $\widetilde{\mathbb{A}}$ have infinite Lebesgue measure and, therefore, a similar proof would fail.

Due to (86) in 3-dimensions the following *conditional probability measures* make sense

$$\lambda_{\widehat{\mathbb{A}}}(\cdot) = \lambda(\cdot | \widehat{\mathbb{A}}) := \frac{\lambda(\cdot \cap \widehat{\mathbb{A}})}{\lambda(\widehat{\mathbb{A}})}, \quad \lambda_{\widetilde{\mathbb{A}}}(\cdot) = \lambda(\cdot | \widetilde{\mathbb{A}}) := \frac{\lambda(\cdot \cap \widetilde{\mathbb{A}})}{\lambda(\widetilde{\mathbb{A}})},$$

and, moreover, due to (85) and (86), for any event $E \in \mathbb{D}$

$$\lim_{r \rightarrow 0} \mu(D_r E | \widehat{\mathbb{A}}_r) = \lambda_{\widehat{\mathbb{A}}}(E), \quad \lim_{r \rightarrow 0} \mu(D_r E | \widetilde{\mathbb{A}}_r) = \lambda_{\widetilde{\mathbb{A}}}(E).$$

In a technical sense, we will only use the upper bound in (85), and (86).

In view of the upper bound in (85), in order to prove (77), (78) and (79) we need, in turn,

$$\lambda \left((u, h, v) \in \widehat{\mathbb{A}} : \angle(-e, \widehat{w}) \leq s \right) \leq C \min\{s, 1\}, \quad (87)$$

$$\lambda \left((u, h, v) \in \widetilde{\mathbb{A}} : \angle(-e, \widetilde{w}) \leq s \right) \leq C \min\{s(|\log s| \vee 1), 1\}, \quad (88)$$

$$\lambda \left((u, h, v) \in \widetilde{\mathbb{A}} : \widetilde{\beta} > s \right) \leq C \min\{s^{-1}(|\log s| \vee 1), 1\}. \quad (89)$$

Here, and in the rest of this section, we use the simplified notation $\widehat{w} := \widehat{w}_1$, $\widetilde{w} := \widetilde{w}_1$, $\widetilde{\beta} := \widetilde{\beta}_1$.

6.5.2. Proof of (87)

Proof. This is straightforward. Recall (75): $\widehat{w}(u, h, v) = u$. For easing notation let

$$\vartheta := \angle(-e, u)$$

and note that for any $t \in \mathbb{R}_+$

$$\left| \{u \in S^2 : 0 \leq \vartheta \leq t\} \right| \leq C \min\{t^2, 1\},$$

with some explicit $C < \infty$.

Then,

$$\begin{aligned} \lambda \left((u, h, v) \in \widehat{\mathbb{A}} : \angle(-e, \widehat{w}) \leq s \right) &\leq \lambda \left((u, h, v) \in \widehat{\mathbb{A}}' : \vartheta \leq s \right) \\ &\leq \lambda \left((u, h, v) \in \mathbb{D} : \vartheta \leq \min\{s, 2h^{-1}\} \right) \\ &= \lambda \left((u, h, v) \in \mathbb{D} : \{h \leq 2s^{-1}\} \cap \{\vartheta \leq s\} \right) \\ &\quad + \lambda \left((u, h, v) \in \mathbb{D} : \{h \geq 2s^{-1}\} \cap \{\vartheta \leq 2h^{-1}\} \right) \\ &\leq Cs. \end{aligned}$$

□

6.5.3. Proof of (88) and (89)

Figure 3 aides understanding this subsection.

Let a and b be the vectors in \mathbb{R}^3 pointing from the origin to the centre of the spherical scatterers of radius 1, on which the first, respectively, the second collision occurs:

$$a = \frac{e - u}{|e - u|}, \quad b = hu + \frac{u - v}{|u - v|},$$

and n the unit vector orthogonal to the plane determined by a and b , pointing so, that $e \cdot n > 0$:

$$n := \frac{a \times b}{|a| |b| \sin(\angle(a, b))},$$

with

$$a \times b = \left(h + \frac{1}{|u - v|}\right) \frac{1}{|e - u|} e \times u - \frac{1}{|e - u| |u - v|} e \times v + \frac{1}{|e - u| |u - v|} u \times v, \quad (90)$$

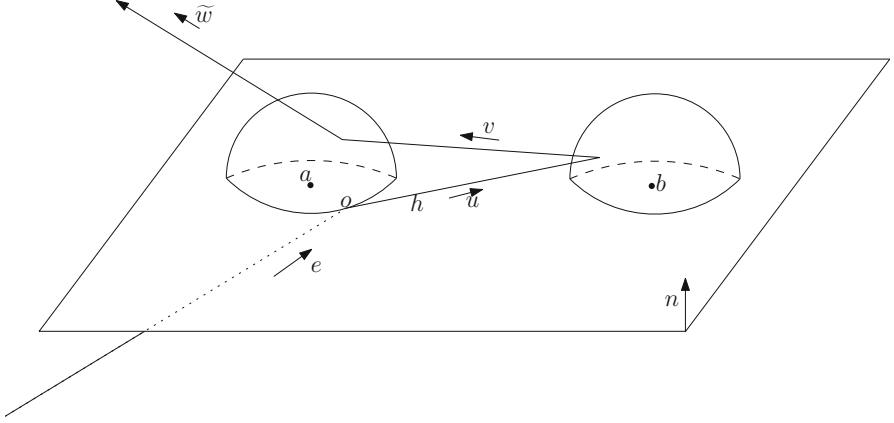


Fig. 3. Above we show a 3 dimensional example of the geometric labelling used in this section. The Z trajectory enters with velocity e from beneath the relevant plane (the dotted line represents motion below the plane). After which the particle remains above the plane

$$|a| = 1, \quad h - 1 \leq |b| \leq h + 1, \quad 0 \leq \sin(\angle(a, b)) \leq 1. \quad (91)$$

Assume there are altogether $\nu \geq 3$ collisions (which occur alternatively, on the first and second scatterer) before escape. Let $w_0 = e$ and w_j , $1 \leq j \leq \nu$, the outgoing velocity after the j -th scattering. So, $w_1 = u$, $w_2 = v$, \dots , $w_\nu = \tilde{w}$.

The proof of (88) and (89) relies on the following observations:

- (a) The n -projection of the velocity of the moving particle does not decrease. More precisely, for $1 \leq j \leq \nu$, $0 \leq w_{j-1} \cdot n \leq w_j \cdot n$. This is due to the choice of the plane determined by the centres of the two scatterers and the first impact point.
- (b) Since $e \cdot n > 0$ and $w_j \cdot n > 0$, for all $1 \leq j \leq \nu$ we have $\angle(-e, w_j) > \frac{\pi}{2} - \angle(n, w_j)$.
- (c) The trapping time $\tilde{\beta}$ is certainly not longer than the time the moving particle spends in the slab $\{x \in \mathbb{R}^3 : 0 \leq x \cdot n \leq 1\}$. Moreover, the scatterers are defocusing, that is, each time there is a collision the velocity component in the n direction increases. Therefore, it follows that

$$\tilde{\beta} \leq h + |v \cdot n|^{-1}. \quad (92)$$

Proof of (88) Without loss of generality we may assume $s \leq \frac{\pi}{2}$.

From the arguments (a) and (b) above it follows, in particular, that

$$\angle(-e, \tilde{w}) = \angle(-e, w_\nu) \geq \frac{\pi}{2} - \angle(n, w_\nu) \geq \frac{\pi}{2} - \angle(n, w_2) = \frac{\pi}{2} - \angle(n, v),$$

and hence

$$\lambda((u, h, v) \in \tilde{\mathbb{A}} : \angle(-e, \tilde{w}) < s) \leq \lambda((u, h, v) \in \tilde{\mathbb{A}}' : |n \cdot v| < 2s). \quad (93)$$

Note that due to (90) and (91)

$$|v \cdot n| \geq \frac{1}{2} |v \cdot (e \times u)|,$$

and thus

$$\lambda \left((u, h, v) \in \widetilde{\mathbb{A}}' : |v \cdot n| < 2s \right) \leq \lambda \left((u, h, v) \in \widetilde{\mathbb{A}}' : |e \cdot (u \times v)| < 4s \right). \quad (94)$$

Next, if u and v are i.i.d. $UNI(S^2)$ -distributed then

$$w := \frac{u \times v}{|u \times v|}, \quad \text{and} \quad \vartheta := |u \times v| = \sin(\angle(u, v))$$

are independent and distributed as

$$w \sim UNI(S^2), \quad \vartheta \sim \mathbb{1}_{\{0 \leq t \leq 1\}}(1 - t^2)^{-1/2} t dt.$$

Therefore,

$$\begin{aligned} & \lambda \left((u, h, v) \in \widetilde{\mathbb{A}}' : |e \cdot (u \times v)| < 4s \right) \\ &= \int_0^\infty dh \int_{S^2} dw \int_0^{\min\{2/h, 1\}} (1 - t^2)^{-1/2} t dt \mathbb{1}_{\{|e \cdot w| \leq \frac{4s}{t}\}} \\ &= \int_0^\infty dh \int_0^{\min\{2/h, 1\}} (1 - t^2)^{-1/2} dt \min\{4s, t\} \\ &\leq C \min\{s |\log s| \vee 1, 1\}. \end{aligned} \quad (95)$$

The last step follows from explicit computations which we omit.

Finally, (93), (94) and (95) yield (88). \square

Proof of (89)

We proceed with the bound (92):

$$\begin{aligned} & \lambda \left((u, h, v) \in \widetilde{\mathbb{A}} : \widetilde{\beta} > s \right) \leq \lambda \left((u, h, v) \in \widetilde{\mathbb{A}}' : h > \frac{s}{2} \right) \\ &+ \lambda \left((u, h, v) \in \widetilde{\mathbb{A}}' : |v \cdot n| < \frac{2}{s} \right). \end{aligned} \quad (96)$$

Bounding the first term on the right hand side of (96) is straightforward:

$$\begin{aligned} & \lambda \left((u, h, v) \in \widetilde{\mathbb{A}}' : h > \frac{s}{2} \right) = \int_{s/2}^\infty \left| \{(u, v) \in S^2 \times S^2 : \angle(-u, v) < 2h^{-1}\} \right| dh \\ &\leq C \int_{s/2}^\infty \min\{h^{-2}, 1\} dh \leq C \min\{s^{-1}, 1\}. \end{aligned} \quad (97)$$

Concerning the second term on the right hand side of (96), this has exactly been done in the proof of (88) above, ending in (95)—with the rôle of s and s^{-1} swapped.

(96), (97) and (95) yield (79). \square

6.6. *Geometric estimates ctd: Proof of Corollary 2.* We start with the following straightforward geometric fact.

Lemma 8. *Let $e, w \in S^2$ and $x \in \mathbb{R}^3$. Then*

$$\left| \{t' > 0 : \min_{t \geq 0} |x + t'w + te| < s\} \right| = \left| \{t' > 0 : \min_{t \geq 0} |x + tw + t'e| < s\} \right| \leq \frac{4s}{\angle(-e, w)}. \quad (98)$$

Proof of Lemma 8. This is elementary 3-dimensional geometry. We omit the details. \square

Proof of (80) and (81) On $\{\widehat{\eta}_k = 1\}$

$$\begin{aligned} \min_{\tau_{k-2} \leq t \leq \tau_k} |Z^{(k)}(t) - Z^{(k)}(\tau_{k-3})| &\geq \min_{0 \leq t} |tu_{k-1} + \xi_{k-2}u_{k-2}| \\ \min_{\tau_{k-3} \leq t \leq \tau_{k-1}} |Z^{(k)}(t) - Z^{(k)}(\tau_k)| &\geq \min_{0 \leq t} \{ \min |\xi_{k-1}u_{k-1} + tu_{k-2} + \xi_k u_{k-1}|, \xi_k \}. \end{aligned} \quad (99)$$

The bounds in (80) and (81) follow from applying (98) and (77), bearing in mind that the distribution density of ξ_{k-2} and ξ_k is bounded. Since these are very similar we will only prove (80) here.

$$\begin{aligned} &\mathbf{P}(\{\widehat{\eta}_k = 1\} \cap \{ \min_{\tau_{k-2} \leq t \leq \tau_k} |Z^{(k)}(t) - Z^{(k)}(\tau_{k-3})| < s \}) \\ &\leq \mathbf{P}(\{\widehat{\eta}_k = 1\} \cap \{ \min_{t \geq 0} |tu_{k-1} + \xi_{k-2}u_{k-2}| < s \}) \\ &= \int_{\widehat{\mathbb{A}}_r} \mathbf{P}(\xi_- \in \{t' : \min_{t \geq 0} |tu + t'e| < s\}) d\mu(u, h, v) \\ &\leq C \int_{\widehat{\mathbb{A}}_r} \min\left\{ \frac{s}{\angle(-e, u)}, 1 \right\} d\mu(u, h, v) \\ &\leq C r s (\log s \vee 1). \end{aligned}$$

In the first step we used (99). The second step follows from the representation (76). The third step relies on (98) and on uniform boundedness of the distribution density of ξ_- (which is either $EXP(1|1)$ or $EXP(1|0)$, depending on the value of ϵ_{k-2}). Finally, the last calculation is based on (77). \square

Proof of (82)

$$\begin{aligned} &\min_{\tau_{k-2} \leq t \leq \tau_k} |Z^{(k)}(t) - Z^{(k)}(\tau_{k-3})| \\ &= \min \left\{ \min_{\tau_{k-2} \leq t \leq \tau_{k-1} + \widetilde{\beta}} |Z^{(k)}(t) - Z^{(k)}(\tau_{k-3})|, \min_{\tau_{k-1} + \widetilde{\beta} \leq t \leq \tau_k} |Z^{(k)}(t) - Z^{(k)}(\tau_{k-3})| \right\}. \end{aligned} \quad (100)$$

Here, and in the rest of this proof, $\widetilde{\beta}$ and \widetilde{w} denote the trapping time and escape direction of the recollision sequence:

$$\widetilde{\beta} := \max\{s \leq \xi_k : \dot{Z}^{(k)}(\tau_{k-1} + s^-) \neq \dot{Z}^{(k)}(\tau_{k-1} + s^+)\} \quad \widetilde{w} := \dot{Z}^{(k)}(\tau_{k-1} + \widetilde{\beta}^+).$$

To bound the first expression on the right hand side of (100) we first observe that by the triangle inequality

$$\min_{\tau_{k-2} \leq t \leq \tau_{k-1} + \tilde{\beta}} \left| Z^{(k)}(t) - Z^{(k)}(\tau_{k-3}) \right| \geq \xi_{k-2} - \xi_{k-1} - 4r. \quad (101)$$

Applying the representation and bounds developed in Sects. 6.4, 6.5,

$$\begin{aligned} & \mathbf{P}(\{\hat{\eta}_k = 0\} \cap \{\tilde{\eta}_k = 1\} \cap \{\min_{\tau_{k-2} \leq t \leq \tau_{k-1} + \tilde{\beta}} |Z^{(k)}(t) - Z^{(k)}(\tau_{k-3})| < s\}) \\ & \leq \mathbf{P}(\{\hat{\eta}_k = 0\} \cap \{\tilde{\eta}_k = 1\} \cap \{\xi_{k-2} \leq \xi_{k-1} + 4r + s\}) \\ & = \int_{\tilde{\mathbb{A}}_r} \mathbf{P}(\xi_- < h + 4r + s) d\mu(u, h, v) \\ & \leq C \int_{\tilde{\mathbb{A}}_r} (\min\{h, 1\} + 4r + s) d\mu(u, h, v) \\ & \leq Cr^2 + Crs + Cr^2 |\log r|. \end{aligned} \quad (102)$$

In the first step we used (101). The second step follows from the representation (76). The third step relies on uniform boundedness of the distribution density of ξ_- (which is either $EXP(1|1)$ or $EXP(1|0)$, depending on the value of ϵ_{k-2}). Finally, the last step follows from explicit calculation, using (85).

To bound the second term on the right hand side of (100) we proceed as in the proof of (80) above. First note that

$$\min_{\tau_{k-1} + \tilde{\beta} \leq t \leq \tau_k} \left| Z^{(k)}(t) - Z^{(k)}(\tau_{k-3}) \right| \geq \min_{0 \leq t} \left| (Z^{(k)}(\tau_{k-2}) - Z^{(k)}(\tau_{k-1} + \tilde{\beta})) + t\tilde{w} + \xi_{k-2}u_{k-2} \right|. \quad (103)$$

Using in turn (103), (76), (98) and uniform boundedness of the distribution density of ξ_- (which is either $EXP(1|1)$ or $EXP(1|0)$, depending on the value of ϵ_{k-2}), and finally (78), we obtain:

$$\begin{aligned} & \mathbf{P}(\{\hat{\eta}_k = 0\} \cap \{\tilde{\eta}_k = 1\} \cap \{\min_{\tau_{k-1} + \tilde{\beta} \leq t \leq \tau_k} |Z^{(k)}(t) - Z^{(k)}(\tau_{k-3})| < s\}) \\ & \leq \mathbf{P}(\{\hat{\eta}_k = 0\} \cap \{\tilde{\eta}_k = 1\} \cap \{\min_{0 \leq t} |(Z^{(k)}(\tau_{k-2}) - Z^{(k)}(\tau_{k-1} + \tilde{\beta})) + t\tilde{w} + \xi_{k-2}u_{k-2}| < s\}) \\ & = \int_{\tilde{\mathbb{A}}_r} \mathbf{P}(\xi_- \in \{t' : \min_{0 \leq t} |\tilde{Z}_r(\tilde{\beta}_r) + t\tilde{w}_r + t'e| < s\}) d\mu(u, h, v) \\ & \leq C \int_{\tilde{\mathbb{A}}_r} \min\{\frac{s}{\angle(-e, \tilde{w}_r)}, 1\} d\mu(u, h, v) \\ & \leq Crs \left(\int_s^\infty \frac{\log(x)}{x} dx \vee 1 \right) \\ & \leq Crs(|\log s|^2 \vee 1). \end{aligned} \quad (104)$$

In the second last line we use (78) and integrate by parts. From (100), (102) and (104) we obtain (82). \square

Proof of (83) We proceed very similarly as in the proof of (82).

$$\min_{\tau_{k-3} \leq t \leq \tau_{k-1} + \tilde{\beta}} \left| Z^{(k)}(t) - Z^{(k)}(\tau_k) \right|$$

$$\geq \min \left\{ \min_{\tau_{k-2} \leq t \leq \tau_{k-1} + \tilde{\beta}} \left| Z^{(k)}(t) - Z^{(k)}(\tau_k) \right|, \min_{\tau_{k-3} \leq t \leq \tau_{k-2}} \left| Z^{(k)}(t) - Z^{(k)}(\tau_k) \right| \right\}. \quad (105)$$

To bound the first expression on the right hand side of (105) we first observe that by the triangle inequality

$$\min_{\tau_{k-2} \leq t \leq \tau_{k-1} + \tilde{\beta}} \left| Z^{(k)}(t) - Z^{(k)}(\tau_k) \right| \geq \xi_k - 2\tilde{\beta} - 4r. \quad (106)$$

Using in turn (106), (76), (79) and explicit computation based on uniform boundedness of the distribution density of ξ_+ (which is either $EXP(1|1)$ or $EXP(1|0)$, depending on the value of ϵ_k) we write

$$\begin{aligned} & \mathbf{P}(\{\widehat{\eta}_k = 0\} \cap \{\widetilde{\eta}_k = 1\} \cap \{ \min_{\tau_{k-2} \leq t \leq \tau_{k-1} + \tilde{\beta}} \left| Z^{(k)}(t) - Z^{(k)}(\tau_k) \right| < s \}) \\ & \leq \mathbf{P}(\{\widehat{\eta}_k = 0\} \cap \{\widetilde{\eta}_k = 1\} \cap \{\xi_k < 8r + 2s\}) + \mathbf{P}(\{\widehat{\eta}_k = 0\} \cap \{\widetilde{\eta}_k = 1\} \cap \{\xi_k < 4\tilde{\beta}\}) \\ & = \mathbf{P}(\xi_+ < 8r + 2s) \mu(\widetilde{\mathbb{A}}_r) + \mathbf{E}(\mu((u, h, v) \in \widetilde{\mathbb{A}}_r : \xi_+ \leq 4\tilde{\beta}_r)) \\ & \leq Cr(r + s) + Cr \mathbf{E}(\min\left\{\left(\frac{\xi_+}{2r}\right)^{-1} \left(\left|\log \frac{\xi_+}{2r}\right| \vee 1\right), 1\right\}) \\ & \leq Cr^2 + Crs + Cr^2 |\log r|^2. \end{aligned} \quad (107)$$

The second term on the right hand side of (105) is bounded in a very similar way as the analogous second term on the right hand side of (100), see (103)-(104). Without repeating these details we state that

$$\mathbf{P}(\{\widehat{\eta}_k = 0\} \cap \{\widetilde{\eta}_k = 1\} \cap \{ \min_{\tau_{k-3} \leq t \leq \tau_{k-2}} \left| Z^{(k)}(t) - Z^{(k)}(\tau_k) \right| < s \}) \leq Crs |\log s|^2. \quad (108)$$

Eventually, from (105), (107) and (108) we obtain (83). \square

6.7. Proof of (66)—concluded. Recall the events $A_{a,b}^{(k)}$, $a, b \in \{1, 2, 3\}$ from the end of Sect. 6.3.

The bounds (80), (81), respectively, (82), (83), with $s = r$, directly imply

$$\begin{aligned} \mathbf{P}(\{\widehat{\eta}_k = 1\} \cap A_{2,2}^{(k)} \mid \underline{\epsilon}) & \leq C\gamma r^2 |\log r|, \\ \mathbf{P}(\{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\} \cap A_{2,2}^{(k)} \mid \underline{\epsilon}) & \leq C\gamma r^2 |\log r|^2. \end{aligned} \quad (109)$$

It remains to prove

$$\begin{aligned} \mathbf{P}(\{\widehat{\eta}_k = 1\} \cap A_{b,2}^{(k)} \mid \underline{\epsilon}) & \leq C\gamma r^2 |\log r|, \\ \mathbf{P}(\{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\} \cap A_{b,2}^{(k)} \mid \underline{\epsilon}) & \leq C\gamma r^2 |\log r|^2, \end{aligned} \quad b = 1, 3. \quad (110)$$

Since the cases $b = 1$ and $b = 3$ are formally identical we will go through the steps of proof with $b = 3$ only. In order to do this we first define the necessary occupation

time measures (Green's functions). For $A \subset \mathbb{R}^3$, define the following occupation time measures for the last part of (71)

$$\begin{aligned}
 G_{\underline{\epsilon}}^{(k)}(A) &:= \mathbf{E}(\#\{1 \leq j \leq \gamma - k : Y(\tau_j) \in A\} \mid \epsilon_{k+j} : 1 \leq j \leq \gamma - k) \\
 &= \mathbf{E}(\#\{k+1 \leq j \leq \gamma : Z^{(k)}(\tau_j) - Z^{(k)}(\tau_k) \in A\} \mid \underline{\epsilon} \cap \{\widehat{\eta}_k = 1\}) \\
 &= \mathbf{E}(\#\{k+1 \leq j \leq \gamma : Z^{(k)}(\tau_j) - Z^{(k)}(\tau_k) \in A\} \mid \underline{\epsilon} \cap \{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\}), \\
 H_{\underline{\epsilon}}^{(k)}(A) &:= \mathbf{E}(\#\{0 \leq t \leq \tau_{\gamma-k} : Y(t) \in A\} \mid \epsilon_{k+j} : 1 \leq j \leq \gamma - k) \\
 &= \mathbf{E}(\#\{\tau_k \leq t \leq \theta : Z^{(k)}(t) - Z^{(k)}(\tau_k) \in A\} \mid \underline{\epsilon} \cap \{\widehat{\eta}_k = 1\}) \\
 &= \mathbf{E}(\#\{\tau_k \leq t \leq \theta : Z^{(k)}(t) - Z^{(k)}(\tau_k) \in A\} \mid \underline{\epsilon} \cap \{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\}).
 \end{aligned}$$

Similarly, define the following occupation time measures for the middle part of (71)

$$\begin{aligned}
 \widehat{G}_{\underline{\epsilon}}^{(k)}(A) &:= \mathbf{E}(\#\{1 \leq j \leq 3 : Z^{(k)}(\tau_{k-j}) - Z^{(k)}(\tau_k) \in A\} \cdot \widehat{\eta}_k \mid \underline{\epsilon}) \\
 \widehat{H}_{\underline{\epsilon}}^{(k)}(A) &:= \mathbf{E}(\#\{\tau_{k-3} \leq t \leq \tau_k : Z^{(k)}(t) - Z^{(k)}(\tau_k) \in A\} \mid \widehat{\eta}_k \mid \underline{\epsilon}) \\
 \widetilde{G}_{\underline{\epsilon}}^{(k)}(A) &:= \mathbf{E}(\#\{1 \leq j \leq 3 : Z^{(k)}(\tau_{k-j}) - Z^{(k)}(\tau_k) \in A\} \cdot \widetilde{\eta}_k \cdot (1 - \widehat{\eta}_k) \mid \underline{\epsilon}) \\
 \widetilde{H}_{\underline{\epsilon}}^{(k)}(A) &:= \mathbf{E}(\#\{\tau_{k-3} \leq t \leq \tau_k : Z^{(k)}(t) - Z^{(k)}(\tau_k) \in A\} \mid \widetilde{\eta}_k \cdot (1 - \widehat{\eta}_k) \mid \underline{\epsilon}).
 \end{aligned}$$

Using the independence of the middle and last parts in the decomposition (71), similarly as (28) or (53), following bounds are obtained

$$\begin{aligned}
 \mathbf{P}(\{\widehat{\eta}_k = 1\} \cap A_{3,2}^{(k)} \mid \underline{\epsilon}) &\leq Cr^{-1} \int_{\mathbb{R}^3} G_{\underline{\epsilon}}^{(k)}(B_{x,2r}) \widehat{H}_{\underline{\epsilon}}^{(k)}(dx) \\
 &\quad + Cr^{-1} \int_{\mathbb{R}^3} H_{\underline{\epsilon}}^{(k)}(B_{x,3r}) \widehat{G}_{\underline{\epsilon}}^{(k)}(dx), \\
 \mathbf{P}(\{\widetilde{\eta}_k = 1\} \cap \{\widehat{\eta}_k = 0\} \cap A_{3,2}^{(k)} \mid \underline{\epsilon}) \\
 &\leq Cr^{-1} \int_{\mathbb{R}^3} G_{\underline{\epsilon}}^{(k)}(B_{x,2r}) \widetilde{H}_{\underline{\epsilon}}^{(k)}(dx) + Cr^{-1} \int_{\mathbb{R}^3} H_{\underline{\epsilon}}^{(k)}(B_{x,3r}) \widetilde{G}_{\underline{\epsilon}}^{(k)}(dx).
 \end{aligned} \tag{111}$$

Due to (34) of Lemma 3 by direct computations the following upper bounds hold

$$G_{\underline{\epsilon}}^{(k)}(B_{x,2r}) \leq CF(|x|), \quad H_{\underline{\epsilon}}^{(k)}(B_{x,3r}) \leq CF(|x|), \tag{112}$$

where $C < \infty$ is an appropriately chosen constant and $F : \mathbb{R}_+ \rightarrow \mathbb{R}$,

$$F(u) := r \mathbb{1}\{0 \leq u < r\} + \frac{r^3}{u^2} \mathbb{1}\{r \leq u < 1\} + \frac{r^3}{u} \mathbb{1}\{1 \leq u < \infty\}.$$

On the other hand, from (80), (81), (82), (83) of Corollary 2 follows that

$$\begin{aligned}
 \widehat{G}_{\underline{\epsilon}}^{(k)}(B_{0,s}) &\leq Crs(|\log s| \vee 1), & \widehat{H}_{\underline{\epsilon}}^{(k)}(B_{0,s}) &\leq Crs(|\log s| \vee 1), \\
 \widetilde{G}_{\underline{\epsilon}}^{(k)}(B_{0,s}) &\leq Cr \max\{s |\log s|^2, r |\log r|^2\} & \widetilde{H}_{\underline{\epsilon}}^{(k)}(B_{0,s}) &\leq Cr \max\{s |\log s|^2, r |\log r|^2\}.
 \end{aligned} \tag{113}$$

Finally, we also have the global bounds

$$\begin{aligned}\widehat{G}_{\underline{\epsilon}}^{(k)}(\mathbb{R}^3) &= 3\mathbf{E}(\widehat{\eta}_k \mid \underline{\epsilon}) \leq Cr, & \widehat{H}_{\underline{\epsilon}}^{(k)}(\mathbb{R}^3) &= \mathbf{E}(\widehat{\eta}_k \cdot \sum_{j=k-2}^k \xi_j \mid \underline{\epsilon}) \leq Cr, \\ \widetilde{G}_{\underline{\epsilon}}^{(k)}(\mathbb{R}^3) &= 3\mathbf{E}(\widetilde{\eta}_k \cdot (1 - \widehat{\eta}_k) \mid \underline{\epsilon}) \leq Cr, & \widetilde{H}_{\underline{\epsilon}}^{(k)}(\mathbb{R}^3) &= \mathbf{E}(\widetilde{\eta}_k \cdot (1 - \widehat{\eta}_k) \cdot \sum_{j=k-2}^k \xi_j \mid \underline{\epsilon}) \leq Cr.\end{aligned}\tag{114}$$

We will prove the upper bound (110) for the first term on the right hand side of the first line in (111). The other four terms are done in very similar way.

First we split the integral as

$$\int_{\mathbb{R}^3} G_{\underline{\epsilon}}^{(k)}(B_{x,2r}) \widehat{H}_{\underline{\epsilon}}^{(k)}(dx) = \int_{|x|<1} G_{\underline{\epsilon}}^{(k)}(B_{x,2r}) \widehat{H}_{\underline{\epsilon}}^{(k)}(dx) + \int_{|x|\geq 1} G_{\underline{\epsilon}}^{(k)}(B_{x,2r}) \widehat{H}_{\underline{\epsilon}}^{(k)}(dx)\tag{115}$$

and note that due to (112) and (114) the second term on the right hand side is bounded as

$$\int_{|x|\geq 1} G_{\underline{\epsilon}}^{(k)}(B_{x,2r}) \widehat{H}_{\underline{\epsilon}}^{(k)}(dx) \leq Cr^4.\tag{116}$$

To bound the first term on the right hand side of (115) we proceed as follows

$$\begin{aligned}\int_{|x|<1} G_{\underline{\epsilon}}^{(k)}(B_{x,2r}) \widehat{H}_{\underline{\epsilon}}^{(k)}(dx) &\leq C \int_0^1 F(u) d\widehat{H}_{\underline{\epsilon}}^{(k)}(B_{0,u}) \\ &= Cr^3 \widehat{H}_{\underline{\epsilon}}^{(k)}(B_{0,1}) - C \int_0^1 \widehat{H}_{\underline{\epsilon}}^{(k)}(B_{0,u}) F'(u) du \\ &\leq Cr^4 + Cr^4 \int_r^1 u^{-2} |\log u| du \\ &\leq Cr^4 + Cr^3 |\log r|.\end{aligned}\tag{117}$$

In the first step we have used (112). The second step is an integration by parts. In the third step we use (113), (114) and the explicit form of the function F . The last step is explicit integration.

Finally, (115), (116), (117) and identical computations for the second term on the right hand side of the first line in (111) yield the first inequality in (110). The second line of (110) for $b = 3$ is proved in an identical way, which we omit to repeat. The case $b = 1$ is done in a formally identical way.

Finally, (66) follows from (73), (109) and (110). \square

7. Proof of Theorem 2—concluded

As in Sect. 4.3 let $\varpi_n = (\gamma_n; (\xi_{n,j}, u_{n,j}) : 1 \leq j \leq \gamma_n)$, $n \geq 1$, be a sequence of i.i.d packs. Denote $\theta_n, ((Y_n(t), Z_n(t)) : 0 \leq t \leq \theta_n)$ the pair of Y and (forward) Z -processes constructed from them and

$$Y(t) = \sum_{k=1}^{v_t} Y(\theta_k) + Y_{v_t+1}(\{t\}), \quad Z(t) = \sum_{k=1}^{v_t} Z(\theta_k) + Z_{v_t+1}(\{t\}).$$

Beside these two we now define yet another auxiliary process $t \mapsto \mathcal{X}(t)$ as follows: $(\mathcal{X}_n(t) : 0 \leq t \leq \theta_n)$ is the Lorentz exploration process constructed with data from $(Y_n(t) : 0 \leq t \leq \theta_n)$ and incoming velocity

$$u_{n,0} = \begin{cases} u_0 & \text{if } n = 1, \\ \dot{\mathcal{X}}_{n-1}(\theta_{n-1}^-) & \text{if } n > 1. \end{cases}$$

Finally, from these legs concatenate

$$\mathcal{X}(t) = \sum_{k=1}^{v_t} \mathcal{X}(\theta_k) + \mathcal{X}_{v_t+1}(\{t\}).$$

Note that the auxiliary process $(\mathcal{X}(t) : 0 \leq t < \infty)$ is not identical with the Lorentz exploration process $(X(t) : 0 \leq t < \infty)$, constructed with data from $(Y(t) : 0 \leq t \leq \infty)$ and initial incoming velocity u_0 , since the former one does not take into account memory effects caused by earlier legs. However, based on Propositions 1 and 2, we will prove that until time $T = T(r) = o(r^{-2} |\log r|^{-2})$ the processes $t \mapsto X(t)$, $t \mapsto \mathcal{X}(t)$, and $t \mapsto Z(t)$ coincide with high probability.

For this, we define the (discrete) stopping times

$$\begin{aligned} \rho &:= \min\{n : \mathcal{X}_n(t) \not\equiv Z_n(t), 0 \leq t \leq \theta_n\} \\ \sigma &:= \min\{n : \max\{\mathbb{1}_{\tilde{W}_n}, \mathbb{1}_{\hat{W}_n} > 0\} = 1\}, \end{aligned}$$

and note that by construction

$$\inf\{t : Z(t) \neq X(t)\} \geq \Theta_{\min\{\rho, \sigma\}-1}.$$

Lemma 9. *Let $T = T(r)$ such that $\lim_{r \rightarrow \infty} T(r) = \infty$ and $\lim_{r \rightarrow \infty} r^2 |\log r|^2 T(r) = 0$. Then*

$$\lim_{r \rightarrow 0} \mathbf{P}(\Theta_{\min\{\rho, \sigma\}-1} < T) = 0. \quad (118)$$

Lemma 10. *Let $T = T(r)$ such that $\lim_{r \rightarrow \infty} T(r) = \infty$ and $\lim_{r \rightarrow \infty} r^2 T(r) = 0$. Then for any $\delta > 0$*

$$\lim_{r \rightarrow 0} \mathbf{P}\left(\max_{0 \leq t \leq T} |Y(t) - Z(t)| > \delta \sqrt{T}\right) = 0. \quad (119)$$

Remark: Actually, (119) holds under the much weaker condition $\lim_{r \rightarrow \infty} r \log \log T = 0$. This can be achieved by applying the LIL rather than a WLLN type of argument to bound $\max_{0 \leq t \leq T} |Y(t) - Z(t)|$ in the proof of Lemma 10, below. However, since the condition of Lemma 9 can not be much relaxed, in the end we would not gain much with the extra effort.

Proof of Lemma 9.

$$\begin{aligned} \mathbf{P}(\Theta_{\min\{\rho, \sigma\}-1} < T) &\leq \mathbf{P}(\rho \leq 2\mathbf{E}(\theta)^{-1}T) + \mathbf{P}(\sigma \leq 2\mathbf{E}(\theta)^{-1}T) + \mathbf{P}\left(\sum_{j=1}^{2\mathbf{E}(\theta)^{-1}T} \theta_j < T\right) \\ &\leq Cr^2 |\log r|^2 T + Cr^2 T + Ce^{-cT}, \end{aligned} \quad (120)$$

where $C < \infty$ and $c > 0$. The first term on the right hand side of (120) is bounded by union bound and (49) from Proposition 1. Likewise, the second term is bounded by union bound and (51) of Proposition 2. In bounding the third term we use a large deviation upper bound for the sum of independent θ_j -s.

Finally, (118) readily follows from (120). \square

Proof of Lemma 10. Note first that

$$\max_{0 \leq t \leq T} |Y(t) - Z(t)| \leq \sum_{j=1}^{v_T+1} \eta_j \xi_j,$$

with v_T and η_j defined in (21), respectively, (24). Hence,

$$\begin{aligned} \mathbf{P}\left(\max_{0 \leq t \leq T} |Y(t) - Z(t)| > \delta\sqrt{T}\right) &\leq \mathbf{P}\left(\sum_{j=1}^{2T} \eta_j \xi_j > \delta\sqrt{T}\right) + \mathbf{P}(v_T > 2T) \\ &\leq C\delta^{-1}\sqrt{T}r + e^{-cT}, \end{aligned} \quad (121)$$

with $C < \infty$ and $c > 0$. The first term on the right hand side of (121) is bounded by Markov's inequality and the straightforward bound

$$\mathbf{E}(\eta_j \xi_j) \leq Cr.$$

The bound on the second term follows from a straightforward large deviation estimate on $v_T \sim POI(T)$.

Finally, (119) readily follows from (121). \square

(9) is direct consequence of Lemmas 9 and 10 and this concludes the proof of Theorem 2. \square

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